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Naveen Velmurugan, Florent Di Meglio

## ▶ To cite this version:

Naveen Velmurugan, Florent Di Meglio. Transient Modelling of Influx and Observer Implementation for Estimation While Drilling. 2018 IEEE Conference on Decision and Control (CDC), Dec 2018, Miami Beach, United States. pp.2623-2628, 10.1109/CDC.2018.8619526. hal-02493310

# HAL Id: hal-02493310 https://minesparis-psl.hal.science/hal-02493310

Submitted on 4 Mar 2020

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## Transient modelling of influx for estimation while drilling

Naveen Velmurugan and Florent Di Meglio

Abstract—In this paper, we present a model of reservoir pressure dynamics in view of estimating influx during drilling. The distributed nature of the model is shown to have an important impact on the transient behaviour of pressure and flow rate when a liquid influx is present. Then, two observers, designed using a backstepping approach, are used to estimate the distributed reservoir pressure as well as wellbore states. The relevance of the approach is illustrated in simulations.

#### I. INTRODUCTION

As automation of the drilling process is becoming an increasingly important research and development focus, real-time measurements are becoming increasingly available and reliable. These enable feedback control, e.g. in Managed Pressure Drilling (MPD) operations [10], as well as improved monitoring. In this paper, we present a distributed model of reservoir dynamics in view of estimating influx during drilling.

The drilling of an oil well consists in creating a borehole several kilometers deep in the ground. While doing so, a drilling fluid is circulated to pressurize the well and remove rock cuttings, among other purposes. The pressure at the bottom of the well, referred to as Bottom Hole Circulating Pressure (BHCP), needs to be accurately controlled to remain inside of uncertain constraints: below the so-called fracture pressure and, in most operations, above the so-called pore (or reservoir) pressure, which is the containment pressure of the fluids in the reservoir. When the BHCP goes below the pore pressure, an influx of fluid flows from the reservoir into the wellbore.

Whether the influx is wanted, as in Under-Balanced Operations (UBO), or accidental, as in most conventional and MPD operations, it is important to estimate its magnitude. Overestimating the severity of a kick produces unnecessary Non-Productive Time, while underestimating it can have disastrous consequences [4].

Influx estimation is an important topic in the literature. In [9], a Kalman filter is applied to a distributed two-phase flow model of the wellbore with a simplified influx model to estimate reservoir characteristics and unmeasured states. In [8], a Lyapunov observer is applied to a finite-dimensional model of the pressure dynamics inside the wellbore with a static reservoir model, to estimate influx and reservoir parameters such as Productivity Index (PI). In [5], the authors rely on downhole flow and pressure measurements to estimate permeability and PI, again, based on a simplified influx model.

In these contributions, the distributed nature of the pressure dynamics, either in the wellbore or the reservoir is often discarded or approximated. The relation between BHCP and influx is either assumed static [8] or formulated as a scalar ODE [9]. These models take their source in production engineering [3], where distributed reservoir dynamics are approximated based on assumptions that are valid in the context of production, usually that either the influx or the pore pressure has reached a quasi-equilibrium.

In this paper, we argue that, in the context of drilling, one should take into account the distributed and transient nature of the near-wellbore reservoir pressure dynamics to estimate the influx. Our approach is as follows. We focus on a reservoir containing only liquid, for which these effects are already important. We present simulations on a model that couples these with distributed pressure dynamics in the drill string and annulus, to illustrate this point. This is the main contribution of the paper.

Then, we recall a backstepping model-based observer design similar to [6] that enables estimation of the distributed reservoir pressure and influx based on two pressure measurements at the wellbore boundaries. Backstepping is a systematic control and observer design method that enables boundary estimation of parabolic [7], [11] and hyperbolic [2], [1] Partial Differential Equations (PDEs), among others. In our case, the system takes the form of a diffusion equation, coupled at its boundary with a system of two counter-convecting hyperbolic PDEs. We use existing results from the literature to derive separate observers for the two systems.

The paper is organized as follows. In Section II, we present the complete model coupling drill string, annulus and reservoir dynamics. In Section III, we recall the observer design. Section IV contains simulations.

#### II. PROBLEM STATEMENT AND MODELLING

We consider the hydraulic part of a MPD drilling system composed of a drill string and an annulus through which flows a liquid, typically oil- or water-based mud. The open part of the annulus is in contact with a reservoir of uncertain pore pressure, permeability and porosity. This setup is schematically depicted on Figure 1. The liquid is injected at a known, variable flow rate using the mud pump at the surface and it is considered that the corresponding flow rate at the bit,  $q_{bit}(t)$  is known. The topside part of the annulus

This research has been carried out in the HYDRA project, which has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement No 675731.

Naveen Velmurugan and Florent Di Meglio are with the Centre Automatique et Systèmes, MINES ParisTech, PSL Research University, 60 bd St-Michel, 75272 Paris Cedex 06 naveen.velmurugan@mines-paristech.fr



Fig. 1. Schematic of flow paths. For simplicity, the drill string and annulus are represented as a u-tube whereas, in reality, they are concentric.

is assumed to be sealed, which is the main feature of MPD operations, and the outflow is controlled via a choke valve. The main purpose of MPD is to improve the control of the BHCP compared to conventional drilling. Among the factors influencing the BHCP are regular operational transients, e.g. when the mud pump flow rate is decreased to zero during a connection, and unwanted incidents, e.g. when the choke is suddenly plugged, or when an unpredicted influx enters the wellbore from the reservoir. Here, we consider the case where a liquid influx is present, due to a reservoir pressure higher than anticipated. To accurately estimate the size of this influx, we model now the coupling between wellbore and reservoir dynamics.

#### A. Wellbore Model

We make the following simplifying assumptions: the flow is one dimensional, single-phase (liquid), temperature is constant. We denote  $\rho$  the liquid density and  $v_l$  its velocity. The liquid is assumed to follow the following Equation of State

$$p(t,x) = p_0 + (\rho(t,x) - \rho_0)C_l^2$$
(1)

where  $C_l$  is the sound velocity in the liquid, t > 0 the time variable and  $x \in [0, L]$  the spatial variable, where *L* is the total length of the pipe. The constants  $p_0$ ,  $\rho_0$  are assumed to be perfectly known in drilling conditions. Writing mass and momentum conservation laws and assuming that

 $v_l \ll C_l$  yields the following set of Partial Differential Equations (PDE) known as Euler's equations,

$$q_t + Aq_x = S \tag{2}$$

where,  $A = \begin{pmatrix} 0 & 1 \\ C_l^2 & 0 \end{pmatrix}$  and  $S = \begin{pmatrix} 0 \\ -F(q) - G(q) \end{pmatrix}$ , F(q) representing frictional pressure losses while G(q) accounts for gravity effects<sup>1</sup>. In what follows we denote

$$\begin{pmatrix} q_1(t,x)\\ q_2(t,x) \end{pmatrix} = \begin{pmatrix} \rho(t,x)\\ \rho(t,x)\nu(t,x) \end{pmatrix}$$
(3)

#### B. Boundary conditions

At the inlet of the drill string, the flow rate is imposed by the pump. At the inlet of the annulus, the volumetric flow rate is given by

$$q_2(t,0) = \rho_0 \frac{q_{bit}(t) + q_{res}(t)}{A}$$
(4)

where A is the annulus area,  $q_{bit}$  is the flow of drilling fluid through the drill bit, given by a one-way-valve-like equation, and  $q_{res}$  is the flow from the reservoir, which will be discussed in the next section. The outflow is governed by a valve equation

$$Aq_{2}(t,L) = K_{c}Z_{c}(t)\sqrt{\frac{2}{q_{1}(t,L)}\left(p(q_{1}(t,L),L) - p_{dc}\right)}$$
(5)

where the downstream choke pressure  $p_{dc}$  is assumed to be constant and  $K_c$  is the choke constant and the flow area through the choke depends on the choke opening  $Z_c$  which is the control input.

In the next section, we detail the model of the reservoir that enables computation of  $q_{res}$ . We also give reduced models typically used for reservoir characterization.

#### C. Reservoir Model

We now derive the equations describing the pressure dynamics in the near-wellbore region of the reservoir. These correspond to the diffusion of hydrocarbons within porous rock. We then recall simplified models typically used in the literature for influx estimation, before making the case for the need to keep the distributed pressure dynamics.

1) Distributed model: We assume that the reservoir section is homogeneous and contains a single type of fluid, with no vertical diffusion. We then obtain the model from the linearized conservation of mass along with Darcy's law relating flow rate to pressure [3]. We denote  $\varphi(r,t)$  as the pressure within the reservoir at a radial location r. The schematic of a radial section of the reservoir is given in Figure 2. The equations inside the radial domain of the reservoir  $r \in [r_w, r_e]$  read

$$\varphi_t(r,t) = \frac{a}{r} (r\varphi_r(r,t))_r \tag{6}$$

<sup>&</sup>lt;sup>1</sup>Notice that, although we do not consider the drill string here, this model is valid there as well, the sign of gravity changing appropriately



Fig. 2. Schematic of a radial section of the reservoir

where  $a = \kappa / (\mu c_t \phi)$  is the diffusivity constant. The influx into the wellbore (or loss) is obtained using, again, Darcy's law and reads

$$q_{res}(t) = \xi \varphi_r(r_w, t) \tag{7}$$

where,  $\xi = 2\pi \kappa r_w h/\mu$ . Notice that the influx depends on the pressure gradient at the boundary of the reservoir that coincides with the wellbore extent. Besides, the pressure in the reservoir at this boundary is assumed to coincide with the BHCP, which yields

$$\varphi(r_w, t) = p(t, 0) \tag{8}$$

Equations (4),(7) and (8) are key as they summarize the coupling between the wellbore and reservoir. We now briefly review existing simplified reservoir models from the literature and compare them with the full distributed model.

2) Simplified models: In [3], the solution to equation (6) with various boundary conditions corresponding to situations arising in oil production is approximated. These lead to the so-called quasi-steady-state solution and constant terminal rate solution.

*a) Quasi-steady-state solution:* A first approximate solution can be obtained by assuming

• that the flux at the radial extent of the reservoir is null, i.e.

$$\varphi_r(r_e,t) = 0 \tag{9}$$

• that the outflow from the reservoir is constant, i.e.  $\dot{q}_{res} = 0$ .

Denoting  $q_{res}^{qss}$  the corresponding reservoir flow, this leads to the following relation between the value of the pressure at the wellbore boundary, at the radial extent and the flow

$$q_{res}^{qss} = -\frac{2\pi\kappa h}{\mu} \frac{(\phi(r_w, t) - p_e)}{(\ln(r_e/r_w) - 1/2)}$$
(10)

Although valid only for stabilized flow conditions, this relation, referred to as the PI relation, is sometimes used in a transient setting [8] or averaged over time [12].



Fig. 3. Comparison of the three reservoir models

*b)* Constant terminal rate solution: Another approximate solution is obtained by assuming

• that the radial extent of the reservoir is infinitely far away, i.e.  $r_e = \infty$  and pressure is constant there, giving the condition

$$\lim_{r \to +\infty} \varphi(r, t) = p_e \tag{11}$$

• that the flow rate  $q_{res}$  is suddenly changed from zero to a constant value

This yields the following relation between flow rate and pressure

$$q_{ctr}(t) = -\frac{4\pi\kappa h(\varphi(r_w, t) - p_e)}{\mu \left(2S_k + \ln\frac{4\kappa t}{\gamma \phi \mu c_r r_w^2}\right)}$$
(12)

This relation is the most widely used for reservoir characterization [9], [13].

c) Comparison of reservoir models: Figure 3 depicts the influx flow rate resulting from these models and a comparison with the full distributed reservoir model (6) along with boundary conditions (8),(9). The scenario corresponds to the liquid influx resulting from a reduction in BHCP caused by opening the choke. Table I summarizes the parameter values used for the simulation. The simulation shows that the influx is largely overestimated by the approximate models, because the conditions under which they are derived (constant flow rate or near-equilibria) are not satisfied in this transient situation. In the case of the quasi-steady-state solution, this is due to the fact that the distributed reservoir pressure reaches the quasi-steady-state only when the well has been producing for several weeks. This is the time needed for the pressure transient to reach the radial extent of the reservoir as described in [3].

Similarly, the constant terminal rate solution relies on a constant flow rate  $q_{res}$ , which is a condition only satisfied several weeks after the start of the production as mentioned in [3]. This type of solution is considered as a basic tool for wellbore pressure analysis that are helpful for the production engineers where the variation in flow rates are obtained

### TABLE I

Parameter	Definition	Value	Unit
$\varphi(r_w,t)$	Pressure at $r_w$	-	Pa
$p_e$	Pressure at $r_e$	1.7e+07	Pa
r <sub>w</sub>	Wellbore radius	0.1143	т
r <sub>e</sub>	Radial extent of reservoir	100	т
μ	Viscosity	5e-03	Pa.s
φ	Porosity	20	%
k	Permeability	5e-12	$m^{-1}$
h	Reservoir thickness	2	т
$S_k$	Skin factor	0	[-]
γ	Constant	1.781	[-]
$c_t$	Total compressibility	2.32e-09	$Pa^{-1}$

RESERVOIR PARAMETERS

by superposition of numerous constant terminal rate solutions [3]. In the interest of estimating influx flow rates while drilling, we are interested in the time period of *at-most a few hours* during which the transient behavior of the reservoir dominates.

In the next section, we derive an observer to estimate the *distributed* reservoir pressure relying on the measurement of BHCP and an observer to estimate the wellbore states relying on the measurement of surface choke pressure.

#### III. OBSERVER DESIGN

In this section, we highlight the structure of the coupling between the wellbore and the reservoir in view of influx estimation. We then present slight modifications of two existing observer designs for the wellbore [2] and reservoir [6], respectively and show their potential for estimation through simulations.

#### A. Coupling boundary conditions

To highlight the structure of the coupling between reservoir and wellbore, we rewrite the wellbore PDE (2) and the linearized form of the boundary conditions (4)–(8) in the Riemann coordinates  $(u, v)^{\top}$  given by

$$\begin{pmatrix} u(t,x) \\ v(t,x) \end{pmatrix} = \frac{1}{2} \begin{pmatrix} C_l q_1(t,x) + q_2(t,x) \\ -C_l q_1(t,x) + q_2(t,x) \end{pmatrix}$$
(13)

This yields a set of linear hyperbolic PDEs of the following form

$$u_t(t,x) + C_l u_x(t,x) = \sigma^{++}(x)u(t,x) + \sigma^{+-}(x)v(t,x) \quad (14)$$

$$v_t(t,x) - C_l v_x(t,x) = \sigma^{-+}(x)u(t,x) + \sigma^{--}(x)v(t,x) \quad (15)$$

where the  $\sigma^{\cdot,\cdot}$  depend on the operating point around which the equations are linearized. Simiarly, the choke equation (5) takes the following form

$$v(t,L) = ku(t,L) + k_U Z_c(t)$$
(16)

where k,  $k_U$  depend, again, on the equilibirum around which we linearize. At the reservoir boundary, Equations (4),(7) and (8) rewrite as

$$u(t,0) = \frac{1}{C_l} \varphi(r_w, t) + v(t,0)$$
(17)

$$\varphi_r(r_w,t) = \frac{A}{\xi} \left( \frac{1}{\rho_0 C_l} \varphi(r_w,t) + \frac{2}{\rho_0} v(t,0) - \frac{q_{bit}(t)}{A} \right) \quad (18)$$

Under the form (17),(18), the boundary conditions give the structure of the coupling: the downward traveling pressure wave v influences the flux out of the reservoir and reflects into the upward traveling pressure wave. In turn, the reservoir pressure affects the upward traveling wave. This interconnection has a feedback structure that potentially generates instability. It is the case, e.g., when the pressure in the wellbore is significantly decreased by an influx of gas from the reservoir, since a larger pressure differential increases the flux. However, in the case considered here of a liquid influx, instability is unlikely as the density of the oil from the reservoir is very close to that of the drilling fluid. Thus, we assume in what follows that the interconnection is stable and design

- an observer for (14)–(17), assuming φ(r<sub>w</sub>,t) is known, and relying on a topside choke pressure measurement p<sub>c</sub>;
- an observer for (6), assuming that v(t,0) is known, relying on a BHCP measurement.



Fig. 4. Schematic of the combined observer design

This yields the observer structure depicted on Figure 4. We now detail the two designs.

#### B. Reservoir Observer

Considering BHCP as a measurement,  $y = \varphi(r_w, t)$ , and assuming that the wellbore state v(t, 0) is known, we consider the following standard Luenberger-like observer

$$\hat{\phi}_t(r,t) = \frac{a}{r} (r \hat{\phi}_r(r,t))_r + l(r) (y - \hat{\phi}(r_w,t))$$
(19)

$$\hat{\varphi}_r(r_e,t) = 0 \tag{20}$$

$$\hat{\varphi}_{r}(r_{w},t) = \frac{A}{\xi} \left( \frac{1}{\rho_{0}C_{l}} \hat{\varphi}(r_{w},t) + \frac{2}{\rho_{0}} v(t,0) - \frac{q_{bit}(t)}{A} \right) + l_{1} \left( y - \hat{\varphi}(r_{w},t) \right) \quad (21)$$

where the observer gains  $l_1$  and l(r) are obtained using the standard backstepping method [6] that provides

$$l_1 = \frac{A}{\xi \rho_0 C_l} + L(r_w, r_w) - K$$
(22)

$$l(r) = aL_y(r, r_w) - (aK + \lambda(r_w))L(r, r_w)$$
(23)

which are obtained by solving the PDE of gain kernel L(r, y) given by,

$$\lambda(y)L_{y}(r,y) + (C + \lambda'(y))L(r,y) + aL_{rr}(r,y) - aL_{yy}(r,y) + \lambda(r)L_{r}(r,y) = 0 \quad (24)$$

where,  $\lambda(r) = a/r$ . The exponential convergence of the resulting error system is proved by mapping it to an exponentially stable target error system using a Volterra coordinate transform as described in [6]. The design parameters C > 0 and K > 0 in the target system are degrees of freedom and can be used to achieve a trade-off between convergence speed, robustness and sensitivity to noise. This point is highlighted in Section IV

#### C. Wellbore Observer

Here, we slightly modify the design from [2], [1] to account for realistic measurements. Indeed, these designs assume that the boundary value of the Riemann invariant u(t,L) can be measured, which is not the case in practice. Rather, pressure sensors are typically available. Here, we generically assume that the measurement rewrites as a linear combination of the Riemann invariants at the topside boundary, which yields

$$y = c_1 u(t, L) + c_2 v(t, L)$$
(25)

In the case of the pressure measurement, one has  $c_1 = -c_2 = C_l$ . We then design, similarly to [2], [1], a Luenberger-like observer as a copy of the plant plus linear output error injection which writes

$$\hat{u}_{t}(t,x) + C_{l}\hat{u}_{x}(t,x) = \sigma^{++}\hat{u}(t,x) + \sigma^{+-}\hat{v}(t,x) - P^{+}(x)(c_{1}\hat{u}(t,L) + c_{2}\hat{v}(t,L) - y)$$
(26)

$$\hat{v}_{t}(t,x) - C_{l}\hat{v}_{x}(t,x) = \sigma^{-+}\hat{u}(t,x) + \sigma^{--}\hat{v}(t,x) - P^{-}(x)(c_{2}\hat{u}(t,L) + c_{2}\hat{v}(t,L) - y)$$
(27)

with the following boundary condition.

$$\hat{v}(t,L) = k\hat{u}(t,L) + \varepsilon y - \varepsilon (c_1\hat{u}(t,L) + c_2\hat{v}(t,L) - y)$$
(28)

$$\hat{u}(t,0) = \frac{1}{C_l} \varphi(r_w, t) + \hat{v}(t,0)$$
(29)

where we assume that the reservoir boundary value  $\varphi(r_w, t)$  is known. Using the standard backstepping approach, the error system  $(\hat{u} - u, \hat{v} - v)$  is mapped to the following target system

$$\alpha_t(t,x) + C_l \alpha_x(t,x) = 0 \tag{30}$$

$$\beta_t(t,x) - C_l \beta_x(t,x) = 0 \tag{31}$$

with boundary conditions

$$\beta(t,L) = \left(\frac{k - \varepsilon c_1}{1 + \varepsilon c_2}\right) \alpha(t,L)$$
(32)

$$\alpha(t,0) = \beta(t,0) \tag{33}$$



Fig. 5. Mapping of effective reflection coefficient as a function of the output error injection gain at the boundary, for different values of the choke opening.

provided the gains  $P^+(x)$  and  $P^-(x)$  satisfy

$$P^{+}(x) = \frac{C_{l}(k - \varepsilon c_{1})P^{uv}(x, L) - C_{l}(1 + \varepsilon c_{2})P^{uu}(x, L)}{c_{1}(1 + \varepsilon c_{2}) + c_{2}(k - \varepsilon c_{1})}$$
(34)  
$$P^{-}(x) = \frac{C_{l}(k - \varepsilon c_{1})P^{vv}(x, L) - C_{l}(1 + \varepsilon c_{2})P^{vu}(x, L)}{c_{1}(1 + \varepsilon c_{2}) + c_{2}(k - \varepsilon c_{1})}$$
(35)

where  $P^{uu}$ ,  $P^{uv}$ ,  $P^{vu}$ ,  $P^{vv}$  are the backstepping kernels from [2]. Similarly to [1], the parameter  $\varepsilon$  is a degree of freedom. More precisely, we denote

$$\bar{k} = \left(\frac{k - \varepsilon c_1}{1 + \varepsilon c_2}\right) \tag{36}$$

the effective reflection coefficient in the target system, which can be freely assigned using the design parameter  $\varepsilon$ . Imposing  $\bar{k}$  close to zero yields a rapidly convergent observer. However, as indicated by (28), large values of  $\varepsilon$  tend to amplify measurement noise. Importantly, a necessary and sufficient condition for the stability of (30)–(33) is that  $\bar{k}$ satisfy

$$-1 < \bar{k} < 1 \tag{37}$$

Notice that the value of k, that comes from the linearization of the choke equation (5), varies according to the equilibrium that, itself, depends on the choke opening. It is possible to change the value of  $\varepsilon$  according to the choke opening to maintain a constant effective reflection coefficient  $\bar{k}$  throughout the operation, but a time-dependent coefficient induces additional dynamics that may lead to instability. Rather, it is possible to map the value of  $\bar{k}$  throughout the operating range and chose a constant epsilon that ensures that (37) is satisfied for all values of  $Z_c$ . This is the point illustrated on Figure 5.

#### **IV. SIMULATIONS**

In this section, we study the coupled observer designed according to the structure depicted on Figure 4. In practice, the estimate  $\hat{v}(t,0)$  from the wellbore observer is inserted

### TABLE II

WELLBORE PARAMETE	RS
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Parameter	Definition	Value	Unit
$q_{bit}$	Flow rate at bit	0.03	$m^3 s^{-1}$
L	Length of the wellbore	2000	т
K <sub>c</sub>	Choke constant	2.85e-03	[-]
$\bar{Z}_c$	Steady state choke opening	0.95	[-]
$C_l$	Velocity of sound	940.30	$m.s^{-1}$
$p_0$	Reference pressure	1e+05	Pa
ρ <sub>0</sub>	Reference density	780	$kg.m^{-3}$

in (21) and the estimate  $\hat{\varphi}(r_w,t)$  from the reservoir observer is inserted in (18) for the reservoir and wellbore observers respectively. We consider various industry relevant scenarios to study the performance and robustness of this combined observer design. Table II summarizes the simulation parameters.

#### A. Drilling into a high pressure pocket

During drilling, it is common to encounter pockets of hydrocarbons with higher pressure than anticipated. Drilling into such zones will lead to an influx. We present here a case where a section with known reservoir pressure is drilled using MPD and a high pressure pocket is encountered at 5s. The pressure within the pocket is considered to be 50% higher than that of the anticipated. As depicted on Figure 6, the observer successfully estimates the influx, while no influx is detected in the absence of an observer (i.e. in open loop). The near wellbore reservoir pressure distribution is also estimated by the observer as shown on Figure 7. Interestingly, the reservoir pressure is poorly estimated away from the wellbore but this does not affect the influx estimation.



Fig. 6. Influx estimation by the observer in comparison to the plant measurement and in the absence of an observer for case–A.

#### B. In presence of transients

During MPD operation, there are instances when the choke opening is varied rapidly in order to maintain the required BHCP. We present the case here with an over estimation



Fig. 7. Time snapshots of distributed pressure profile within the reservoir obtained from plant in comparison to that of the estimation by the observer for case–A.

of the reservoir and wellbore parameters to emulate the uncertainty of these parameters in practice, inclusive of transients introduced by the choke variations. The perturbation of wellbore states from the steady state profile, reservoir pressure and permeability are assumed to be overestimated by 10% each and  $Z_c$  varied from 0.8 to 0.4 in 1s. Despite the uncertainty in the reservoir and wellbore parameters, the observer successfully estimates the influx and the near wellbore reservoir pressure distribution as depicted in Figures 8 and 9 respectively.



Fig. 8. Influx estimation by the observer in comparison to the plant measurement for case–B.

#### V. CONCLUSIONS AND FUTURE WORKS

We have presented a reservoir model capturing the distributed properties of the pressure dynamics in view of influx estimation. Applying existing observer designs for the wellbore and reservoir dynamics enables estimation of the influx flow rate despite uncertainty on the reservoir characteristics. This first step opens a large perspective of future works.



Fig. 9. Time snapshots of distributed pressure profile within the reservoir obtained from plant in comparison to that of the estimation by the observer for case–B.

First, a coupled parabolic-hyperbolic observer could be designed, which would be a novel result in control and estimation of PDEs. Modifying the backstepping transformations such that the entire wellbore-reservoir state is mapped to a stable target system induces nonlinear, nonlocal relations between the backstepping kernels, which severely complicates the proof of their well-posedness.

Most importantly, influxes of fluids with densities varying from the drilling fluid need to be considered and modeled. This drastically modifies the wellbore dynamics, in particular if gas is present. This is a key direction for future work. Another important modelling aspect is to deal with the potential longitudinal heterogeneity of the reservoir in the open section of the well.

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