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Felix Antritter, Jean Lévine

# Flatness Characterization: Two Approaches 

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#### Abstract

We survey two approaches to flatness necessary and sufficient conditions and compare them on examples.


## 1 Introduction

In this survey we consider underdetermined implicit systems of the form

$$
\begin{equation*}
F(x, \dot{x})=0 \tag{1}
\end{equation*}
$$

with $x \in X, X$ being an inifnitely differentiable manifold of dimension $n$, whose tangent bundle is denoted by $\mathrm{T} X$, and $F: \mathrm{T} X \rightarrow \mathbb{R}^{n-m}$ regular in the sense that rk $\frac{\partial F}{\partial \dot{x}}=n-m$ in a suitable open dense subset of TX. Differential flatness, or more shortly, flatness was introduced in 1992 [20|11]. In the setting of implicit control systems it may be roughly described as follows: there exists a smooth mapping $x=\varphi\left(y, \dot{y}, \ldots, y^{(r)}\right)$ with $y=\left(y_{1}, \ldots, y_{m}\right)^{T}$ of dimension $m, r=\left(r_{1}, \ldots, r_{m}\right)^{T} \in \mathbb{N}^{m}$, such that

$$
\begin{equation*}
F\left(\varphi\left(y, \dot{y}, \ldots, y^{(r)}\right), \dot{\varphi}\left(y, \dot{y}, \ldots, y^{(r+1)}\right)\right) \equiv 0 \tag{2}
\end{equation*}
$$

with $\varphi$ invertible in the sense that there exists a locally defined smooth mapping $\psi$ and a multi-index $s$ such that $y=\psi\left(x, \dot{x}, \ldots, x^{(s)}\right)$.
The vector $y$ is called a flat output.
This concept has inspired an important literature. See [10|21|19|26|27|31] for surveys on flatness and its applications. Various formalisms have been introduced: finite dimensional differential geometric approaches [4|14|30, [32|28, differential algebra and related approaches 12|3|15, infinite dimensional differential geometry of jets and prolongations $13|33| 19|6| 7 \mid 23$, [22|24, which is adopted here. The interested reader may refer to [1|13|16], [19|23|34 for more details.
The first part of the paper recalls the mathematical setting. In Section 3 the approch introduced in [19|2] for the characterization of differentially flat systems is recalled. Then, in Section 4 we introduce a novel characterization using the so-called Generalized Euler-Lagrange Operator. We conclude the paper with examples.

## 2 Implicit control systems on manifolds of jets of infinite order

Given an infinitely differentiable manifold $X$ of dimension $n$, we denote its tangent space at $x \in X$ by $\mathrm{T}_{x} X$, and its tangent bundle by $\mathrm{T} X$. Let $F$ be a meromorphic function from $\mathrm{T} X$ to $\mathbb{R}^{n-m}$. We consider an underdetermined implicit system of the form (1) regular in the sense that rk $\frac{\partial F}{\partial \dot{x}}=n-m$ in a suitable open dense subset of $\mathrm{T} X$.
Following 17|18, we consider the infinite dimensional manifold $\mathfrak{X}$ defined by $\mathfrak{X} \stackrel{\text { def }}{=} X \times \mathbb{R}_{\infty}^{n} \stackrel{\text { def }}{=} X \times \mathbb{R}^{n} \times \mathbb{R}^{n} \times \ldots$, made of an infinite (but countable) number of copies of $\mathbb{R}^{n}$, with the global infinite set of coordinates $s^{3} \bar{x}=\left(x, \dot{x}, \ldots, x^{(k)}, \ldots,\right)$, endowed with the product topology. Recall that, in this topology, a function $\varphi$ from $\mathfrak{X}$ to $\mathbb{R}$ is continuous (resp. differentiable) if $\varphi$ depends only on a finite (but otherwise arbitrary) number of variables and is continuous (resp. differentiable) with respect to these variables. $C^{\infty}$ or analytic or meromorphic functions from $\mathfrak{X}$ to $\mathbb{R}$ are then defined as in the usual finite dimensional case since they only depend on a finite number of variables. We endow $\mathfrak{X}$ with the so-called trivial Cartan field ([16|34]) $\tau_{\mathfrak{X}}=\sum_{i=1}^{n} \sum_{j \geq 0} x_{i}^{(j+1)} \frac{\partial}{\partial x_{i}^{(j)}}$. We also denote by $L_{\tau_{\mathcal{X}}} \gamma=\sum_{i=1}^{n} \sum_{j \geq 0} x_{i}^{(j+1)} \frac{\partial \gamma}{\partial x_{i}^{(j)}}=\frac{d \gamma}{d t}$ the Lie derivative of a differentiable function $\gamma$ along $\tau_{\mathfrak{X}}$ and $L_{\tau_{\mathfrak{X}}}^{k} \gamma$ its $k$ th iterate. Since $\frac{d}{d t} x_{i}^{(j)} \stackrel{\text { def }}{=} \dot{x}_{i}^{(j)}=x_{i}^{(j+1)}$, the Cartan field acts on coordinates as a shift to the right. $\mathfrak{X}$ is thus called manifold of jets of infinite order.
A regular implicit control system is defined as a triple $\left(\mathfrak{X}, \tau_{\mathfrak{X}}, F\right)$ with $\mathfrak{X}=X \times \mathbb{R}_{\infty}^{n}, \tau_{\mathfrak{X}}$ its associated trivial Cartan field, and $F$ meromorphic from $\mathrm{T} X$ to $\mathbb{R}^{n-m}$ ) satisfying $\mathrm{rk} \frac{\partial F}{\partial \dot{x}}=n-m$ in a suitable open subset of $\mathrm{T} X$.
We next consider the cotangent space $T_{\bar{x}}^{*} \mathfrak{X}$ with $d x_{i}^{(j)}, i=1, \ldots, n, j \geq 0$ as basis, dual to the $\frac{\partial}{\partial x_{i}^{(j)}}$ 's. 1-forms on $\mathfrak{X}$ are then defined in the usual way. The set of 1 -forms is noted $\Lambda^{1}(\mathfrak{X})$. We also denote by $\Lambda^{p}(\mathfrak{X})$ the module of all the $p$-forms on $\mathfrak{X}$.

### 2.1 Flatness

We recall the following definitions and result $17|18| 19$ :
Given two regular implicit control systems $\left(\mathfrak{X}, \tau_{\mathfrak{X}}, F\right)$, with $\mathfrak{X}=X \times \mathbb{R}_{\infty}^{n}$, $\operatorname{dim} X=n$ and $\operatorname{rk} \frac{\partial F}{\partial \dot{x}}=n-m$, and $\left(\mathfrak{Y}, \tau_{\mathfrak{V}}, G\right)$, with $\mathfrak{Y}=Y \times \mathbb{R}_{\infty}^{p}$, $\operatorname{dim} Y=p, \tau_{\mathfrak{Y}}$ its trivial Cartan field, and rk $\frac{\partial G}{\partial \dot{y}}=p-q$, we set $\mathfrak{X}_{0}=\{\bar{x} \in$ $\left.\mathfrak{X} \mid L_{\tau_{\mathfrak{X}}}^{k} F(\bar{x})=0, \forall k \geq 0\right\}$ and $\mathfrak{Y}_{0}=\left\{\bar{y} \in \mathfrak{Y} \mid L_{\tau_{\mathfrak{Y}}}^{k} G(\bar{y})=0, \forall k \geq 0\right\}$. They are endowed with the topologies and differentiable structures induced by $\mathfrak{X}$ and $\mathfrak{Y}$ respectively.

Definition 1 The control systems $\left(\mathfrak{X}, \tau_{\mathfrak{X}}, F\right)$ and $\left(\mathfrak{Y}, \tau_{\mathfrak{Y}}, G\right)$ are said locally Lie-Bäcklund equivalent (or shortly L-B equivalent) in a neighbourhood $X_{0} \times y_{0}$ of the pair $\left(\bar{x}_{0}, \bar{y}_{0}\right) \in \mathfrak{X}_{0} \times \mathfrak{Y}_{0}$ if and only if

[^0](i) there exists a one-to-one meromorphic mapping $\Phi=(\varphi, \dot{\varphi}, \ldots)$ from $y_{0}$ to $X_{0}$ satisfying $\Phi\left(\bar{y}_{0}\right)=\bar{x}_{0}$ and such that $\Phi_{*} \tau_{\mathfrak{Y}}=\tau_{\mathfrak{X}}$;
(ii) there exists $\Psi$ one-to-one and meromorphic from $\mathcal{X}_{0}$ to $y_{0}$, with $\Psi=$ $(\psi, \dot{\psi}, \ldots)$, such that $\Psi\left(\bar{x}_{0}\right)=\bar{y}_{0}$ and $\Psi_{*} \tau_{\mathfrak{X}}=\tau_{\mathfrak{Y}}$.
The mappings $\Phi$ and $\Psi$ are called mutually inverse Lie-Bäcklund isomorphisms at $\left(\bar{x}_{0}, \bar{y}_{0}\right)$.

Definition 2 The implicit system $\left(\mathfrak{X}, \tau_{\mathfrak{X}}, F\right)$ is locally flat in a neighborhood of $\left(\bar{x}_{0}, \bar{y}_{0}\right) \in \mathfrak{X}_{0} \times \mathbb{R}_{\infty}^{m}$ if and only if it is locally L-B equivalent around $\left(\bar{x}_{0}, \bar{y}_{0}\right)$ to the trivial implicit system $\left(\mathbb{R}_{\infty}^{m}, \tau_{\mathbb{R}_{\infty}^{m}}^{m}, 0\right)$. In this case, the mutually inverse L-B isomorphisms $\Phi$ and $\Psi$ are called inverse trivializations.

Theorem 1 The system $\left(\mathfrak{X}, \tau_{\mathfrak{X}}, F\right)$ is locally flat at $\left(\bar{x}_{0}, \bar{y}_{0}\right) \in \mathfrak{X}_{0} \times \mathbb{R}_{\infty}^{m}$ if and only if there exists a local meromorphic invertible mapping $\Phi$ from $\mathbb{R}_{\infty}^{m}$ to $\mathfrak{X}_{0}$, with meromorphic inverse, satisfying $\Phi\left(\bar{y}_{0}\right)=\bar{x}_{0}$, and such tha ${ }^{4}$

$$
\begin{equation*}
\Phi^{*} d F=0 . \tag{3}
\end{equation*}
$$

## 3 Necessary and Sufficient Conditions: Generalized Moving Frame Structure Equations

### 3.1 Algebraic characterization of the differential of a trivialization

Consider the following matrix, polynomial with respect to the differential operator $\frac{d}{d t}$ (we use indifferently $\frac{d}{d t}$ for $L_{\tau_{\mathcal{X}}}$ or $L_{\tau_{\mathbb{R}_{\infty}^{m}}}$, the context being unambiguous):

$$
\begin{equation*}
P(F)=\frac{\partial F}{\partial x}+\frac{\partial F}{\partial \dot{x}} \frac{d}{d t}, \quad P(\varphi)=\sum_{j \geq 0} \frac{\partial \varphi}{\partial y^{(j)}} \frac{d^{j}}{d t^{j}} \tag{4}
\end{equation*}
$$

with $P(F)($ resp. $P(\varphi))$ of size $(n-m) \times n($ resp. $n \times m)$.
Equation (3) reads:

$$
\begin{equation*}
\Phi^{*} d F=P(F) P(\varphi) d y=0 . \tag{5}
\end{equation*}
$$

Clearly, the entries of the matrices in (4) are polynomials in the differential operator $\frac{d}{d t}$ with meromorphic coefficients from $\mathfrak{X}$ to $\mathbb{R}$. We denote by $\mathfrak{K}$ the field of meromorphic functions from $\mathfrak{X}$ to $\mathbb{R}$ and by $\mathfrak{K}\left[\frac{d}{d t}\right]$ the (noncommutative) principal ideal ring of polynomials in $\frac{d}{d t}$ with coefficients in $\mathfrak{K}$. For $r, s \in \mathbb{N}$, let us denote by $\mathcal{M}_{r, s}\left[\frac{d}{d t}\right]$ the module of $r \times s$ matrices over $\mathfrak{K}\left[\frac{d}{d t}\right]$ (see e.g. [8]). Matrices whose inverse belong to $\mathcal{M}_{r, r}\left[\frac{d}{d t}\right]$ are called unimodular matrices. They form a multiplicative group denoted by $\mathcal{U}_{r}\left[\frac{d}{d t}\right]$.

[^1]Every matrix $M \in \mathcal{M}_{r, s}\left[\frac{d}{d t}\right]$ admits a Smith decomposition (or diagonal reduction)

$$
\begin{equation*}
V M U=\left(\Delta, 0_{r, s-r}\right) \text { if } r \leq s, \text { and }\binom{\Delta}{0_{r-s, s}} \text { if } s \leq r \tag{6}
\end{equation*}
$$

with $V \in \mathcal{U}_{r}\left[\frac{d}{d t}\right]$ and $U \in \mathcal{U}_{s}\left[\frac{d}{d t}\right]$ and $\Delta$ diagonal (see e.g. [8]). $U$ and $V$ are indeed non unique. We say that $U \in \mathrm{R}-\operatorname{Smith}(M)$ and $V \in$ $\mathrm{L}-\operatorname{Smith}(M)$.
A matrix $M \in \mathcal{M}_{r, s}\left[\frac{d}{d t}\right]$ is said hyper-regular if and only if its Smith decomposition leads to $\Delta=I$. An interpretation of this property in terms of controllability in the sense of (9), may be found in [18.
From now on, we assume that $P(F)$ is hyper-regular in a neighborhood of $\bar{x}_{0}$. In place of (5), we first solve the matrix equation:

$$
\begin{equation*}
P(F) \Theta=0 \tag{7}
\end{equation*}
$$

where $\Theta \in \mathcal{M}_{n, m}\left[\frac{d}{d t}\right]$ is not supposed to be of the form $P(\varphi)$. It may be verified that matrices $\Theta \in \mathcal{M}_{n, m}\left[\frac{d}{d t}\right]$ satisfying $[7]$ have the structure

$$
\begin{equation*}
\Theta=U\binom{0_{n-m, m}}{I_{m}} W \tag{8}
\end{equation*}
$$

with $U \in \mathrm{R}$ - Smith $(P(F))$ and $W \in \mathcal{U}_{m}\left[\frac{d}{d t}\right]$ arbitrary. Clearly $\Theta$ is itself hyper-regular and admits the Smith decomposition

$$
\begin{array}{r}
Q \Theta Z=Q U\binom{0_{n-m, m}}{I_{m}} W Z=Q \hat{U} R=\binom{I_{m}}{0_{n-m, m}}  \tag{9}\\
\text { with } Q \in U_{n}\left[\frac{d}{d t}\right], Z \in U_{m}\left[\frac{d}{d t}\right], R=W Z \text { and } \hat{U}=U\binom{0_{n-m, m}}{I_{m}} .
\end{array}
$$

### 3.2 Integrability

We denote by $\omega$ the $m$-dimensional vector 1 -form defined by

$$
\left.\omega(\bar{x})=\left(\begin{array}{c}
\omega_{1}(\bar{x})  \tag{10}\\
\vdots \\
\omega_{m}(\bar{x})
\end{array}\right)=\left(I_{m}, 0_{m, n-m}\right) Q(\bar{x}) d x \right\rvert\, x_{0}
$$

with $Q$ given by $\sqrt{9}$, the restriction to $\mathfrak{X}_{0}$ meaning that $\bar{x} \in X_{0}$ satisfies $L_{\tau_{\mathfrak{X}}}^{k} F=0$ for all $k$ and that the $d x_{j}^{(k)}$ are such that $d L_{\tau_{\mathcal{X}}}^{k} F=0$ in $X_{0}$ for all $k$. Since $Q$ is hyper-regular, the forms $\omega_{1}, \ldots, \omega_{m}$ are independent by construction.

Theorem 2 A necessary and sufficient condition for system (1) to be locally flat around $\left(\bar{x}_{0}, \bar{y}_{0}\right)$ is that there exist $U \in \mathrm{R}-\operatorname{Smith}(P(F))$, $Q \in \mathrm{~L}-\operatorname{Smith}(\hat{U})$, with $\hat{U}$ given by (9) and a matrix $M \in \mathcal{U}_{m}\left[\frac{d}{d t}\right]$ such that $d(M \tau)=0$.

We denote by $\left(\Lambda^{p}(\mathfrak{X})\right)^{m}$ the space of $m$-dimensional vector $p$-forms on $\mathfrak{X}$, by $(\Lambda(\mathfrak{X}))^{m}$ the space of $m$-dimensional vector forms of arbitrary degree on $\mathfrak{X}$, and by $\mathcal{L}_{q}\left((\Lambda(\mathfrak{X}))^{m}\right)=\bigcup_{p \geq 1} \mathcal{L}\left(\left(\Lambda^{p}(\mathfrak{X})\right)^{m},\left(\Lambda^{p+q}(\mathfrak{X})\right)^{m}\right)$ the space of linear operators from $\left(\Lambda^{p}(\mathfrak{X})\right)^{m}$ to $\left(\Lambda^{p+q}(\mathfrak{X})\right)^{m}$ for all $p \geq 1$, where $\mathcal{L}(\mathcal{P}, \mathcal{Q})$ denotes the set of linear mappings from a given space $\mathcal{P}$ to a given space $\mathcal{Q}$.
In order to develop the expression $d(\mu \kappa)$ for $\mu \in \mathcal{L}_{q}\left((\Lambda(\mathfrak{X}))^{m}\right)$ and for all $\kappa \in\left(\Lambda^{p}(\mathfrak{X})\right)^{m}$ and all $p \geq 1$, we define the operator $\mathfrak{d}$ by:

$$
\begin{equation*}
\mathfrak{d}(\mu) \kappa=d(\mu \kappa)-(-1)^{q} \mu d \kappa . \tag{11}
\end{equation*}
$$

Note that 11 uniquely defines $\mathfrak{d}(\mu)$ as an element of $\mathcal{L}_{q+1}\left((\Lambda(\mathfrak{X}))^{m}\right)$.
Theorem 3 The system $\left(\mathfrak{X}, \tau_{\mathfrak{X}}, F\right)$ is locally flat iff there locally exists $\mu \in \mathcal{L}_{1}\left((\Lambda(\mathfrak{X}))^{m}\right)$, and a matrix $M \in \mathcal{U}_{m}\left[\frac{d}{d t}\right]$ such that

$$
\begin{equation*}
d \omega=\mu \omega, \quad \mathfrak{d}(\mu)=\mu^{2}, \quad \mathfrak{d}(M)=-M \mu . \tag{12}
\end{equation*}
$$

with the notation $\mu^{2}=\mu \mu$ and where $\omega$ is defined by 10. In addition, if (12) holds true, a flat output $y$ is obtained by integration of $d y=M \omega$.

Remark 1 Note that the two first conditions of (12) are comparable to conditions ( $A$ ) and ( $B$ ) of [ 6 (7]. However, the last condition of (12) is different from condition (C) of [6] 7] and is easier to check.
Note also that conditions (12) may be seen as a generalization in the framework of manifolds of jets of infinite order of Cartan's well-known moving frame structure equations (see e.g. [5]).

### 3.3 A Sequential Procedure

We start with $P(F)$ hyper-regular and compute the vector 1-form $\omega$ defined by 10 .
1 . We identify the operator $\mu$ such that $d \omega=\mu \omega$ componentwise. It is proven in 19 that such $\mu$ always exists.
2. Among the possible $\mu$ 's, only those satisfying $\mathfrak{d}(\mu)=\mu^{2}$ are kept. It is shown in 19 that such $\mu$ always exists.
3 . We then identify $M$ such that $\mathfrak{d}(M)=-M \mu$ componentwise.
4. If, among such $M$ 's, there is a unimodular one, the system is flat and a flat output is obtained by integration of $d y=M \omega$. Otherwise the system is not flat.
More details and examples may be found in (18|19.

## 4 Necessary and Sufficient Conditions using the Generalized Euler-Lagrange Operator

Another way of analysing (3) consists in characterizing the change of coordinates corresponding to the mapping $\Phi$ in (3). More precisely (3) reads

$$
\begin{equation*}
\sum_{j=1}^{m} \sum_{k=0}^{r_{j}}\left(\frac{\partial F}{\partial x} \frac{\partial \varphi}{\partial y_{j}^{(k)}} d y_{j}^{(k)}+\frac{\partial F}{\partial \dot{x}} \frac{d}{d t}\left(\frac{\partial \varphi}{\partial y_{j}^{(k)}}\right) d y_{j}^{(k)}+\frac{\partial F}{\partial \dot{x}} \frac{\partial \varphi}{\partial y_{j}^{(k)}} d y_{j}^{(k+1)}\right)=0 \tag{13}
\end{equation*}
$$

Since the one forms $d y_{1}, \ldots, d y_{1}^{\left(r_{1}\right)}, \ldots, d y_{m}, \ldots, d y_{m}^{\left(r_{m}\right)}$ are independent by assumption, 13 yields, for every $j=1, \ldots, m$,

$$
\left\{\begin{array}{l}
\frac{\partial F}{\partial \dot{x}} \frac{\partial \varphi}{\partial y_{j}^{\left(r_{j}\right)}}=0  \tag{14}\\
\frac{\partial F}{\partial x} \frac{\partial \varphi}{\partial y_{j}^{(k)}}+\frac{\partial F}{\partial \dot{x}} \frac{d}{d t}\left(\frac{\partial \varphi}{\partial y_{j}^{(k)}}\right)+\frac{\partial F}{\partial \dot{x}} \frac{\partial \varphi}{\partial y_{j}^{(k-1)}}=0, \quad \forall k=1, \ldots, r_{j} \\
\frac{\partial F}{\partial x} \frac{\partial \varphi}{\partial y_{j}}+\frac{\partial F}{\partial \dot{x}} \frac{d}{d t}\left(\frac{\partial \varphi}{\partial y_{j}}\right)=0
\end{array}\right.
$$

The Generalized Euler-Lagrange operator $\varepsilon_{F}$ associated to $F$ is defined by

$$
\begin{equation*}
\mathcal{E}_{F}=\frac{\partial F}{\partial x}-\frac{d}{d t}\left(\frac{\partial F}{\partial \dot{x}}\right) \tag{15}
\end{equation*}
$$

In the case $n-m=1$, it is well-known that the curves that extremize the cost function $J=\int_{0}^{T} F(x, \dot{x}) d t$ are those satisfying the Euler-Lagrange equation $\mathcal{E}_{F}=0$, which justifies our terminology.
Using (15) and elementary calculus, (14) yields:
Theorem $4 A$ necessary and sufficient condition for (1) to be diffferentially flat is that there exist $\left(r_{1}, \ldots, r_{m}\right)$ with $\sum_{i=1}^{m} r_{i}+m \geq n$ and a solution $\varphi$ of the following triangular system of PDEs in an open dense subset of $\mathfrak{X}$

$$
\left\{\begin{array}{l}
\frac{\partial F}{\partial \dot{x}} \frac{\partial \varphi}{\partial y_{j}^{\left(r_{j}\right)}}=0  \tag{16}\\
\frac{\partial F}{\partial \dot{x}} \frac{\partial \varphi}{\partial y_{j}^{(l)}}=\sum_{k=0}^{r_{j}-l-1}(-1)^{k+1} \frac{d^{k}}{d t^{k}}\left(\varepsilon_{F} \frac{\partial \varphi}{\partial y_{j}^{(l+k+1)}}\right), \quad \forall l=0, \ldots, r_{j}-1 \\
0=\sum_{k=0}^{r_{j}}(-1)^{k} \frac{d^{k}}{d t^{k}}\left(\varepsilon_{F} \frac{\partial \varphi}{\partial y_{j}^{(k)}}\right)
\end{array}\right.
$$

satisfying $d \varphi_{1} \wedge \ldots \wedge d \varphi_{n} \neq 0$.
Remark 2 If there exists a coordinate transformation $\varphi$ that satisfies the conditions of Theorem 4 with given $r_{1}, \ldots, r_{m}$, meaning that the system is flat, then $g_{j}=\sum_{i=1}^{n} \frac{\partial \varphi_{i}}{\left.\partial y_{j} r_{j}\right)} \frac{\partial}{\partial \dot{x}_{i}}$, if non zero, defines a ruled direction [32[25 19].

## 5 Examples

### 5.1 An Academic Example: Generalized Moving Frame Approach

We consider the 3 -dimensional system with 2 inputs:

$$
\begin{equation*}
\dot{x}_{1}=u_{1}, \quad \dot{x}_{2}=u_{2}, \quad \dot{x}_{3}=\sin \left(\frac{u_{1}}{u_{2}}\right) \tag{17}
\end{equation*}
$$

or, in implicit form:

$$
\begin{equation*}
F\left(x_{1}, x_{2}, x_{3}, \dot{x}_{1}, \dot{x}_{2}, \dot{x}_{3}\right) \triangleq \dot{x}_{3}-\sin \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right)=0 \tag{18}
\end{equation*}
$$

It is readily seen that $P(F)=\left[-\cos \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right) \dot{x}_{2}^{-1} \frac{d}{d t}\left|\dot{x}_{1} \cos \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right) \dot{x}_{2}^{-2} \frac{d}{d t}\right| \frac{d}{d t}\right]$ and that $V P(F) U=\left(\begin{array}{lll}1 & 0 & 0\end{array}\right)$ with

$$
V=1, \quad U=\left(\begin{array}{c|c|c}
\frac{\dot{x}_{1}}{a \dot{x}_{2}} & 1+\frac{\dot{x}_{1}}{a\left(\dot{x}_{2}\right)^{2}} \cos \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right) \frac{d}{d t} & \frac{\dot{x}_{1}}{a \dot{x}_{2}} \frac{d}{d t}  \tag{19}\\
\frac{1}{a} & \frac{1}{a \dot{x}_{2}} \cos \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right) \frac{d}{d t} & -\frac{1}{a} \frac{d}{d t} \\
0 & 0 & 1
\end{array}\right)
$$

where $a=-\frac{1}{\dot{x}_{2}} \cos \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right)\left(\frac{\ddot{x}_{1} \dot{x}_{2}-\dot{x}_{1} \dot{x}_{2}}{\left(\dot{x}_{2}\right)^{2}}\right)$. Then, $Q \hat{U} R=\left(\begin{array}{ll}1 & 0 \\ 0 & 1 \\ 0 & 0\end{array}\right)$ is computed with

$$
Q=\left(\begin{array}{c|c|c}
1 & -\frac{\dot{x}_{1}}{\dot{x}_{2}} & 0  \tag{20}\\
0 & 0 & 1 \\
-\frac{1}{a \dot{x}_{2}} \cos \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right) \frac{d}{d t} & \frac{\dot{x}_{1}}{a\left(\dot{x}_{2}\right)^{2}} \cos \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right) & \frac{d}{d t}
\end{array} \frac{1}{a} \frac{d}{d t}\right), \quad R=\left(\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right)
$$

So, $\quad\left(\omega_{1} \omega_{2}\right)^{T}=\left(\begin{array}{lll}1 & 0 & 0 \\ 0 & 0 & 1\end{array}\right) Q d x=\left(d x_{1}-\frac{\dot{x}_{1}}{\dot{x}_{2}} d x_{2} d x_{3}\right)^{T}$ and $d \omega=\left(\frac{1}{\sqrt{1-\left(x_{3}\right)^{2}}} d x_{2} \wedge d x_{3} 0\right)^{T}$. According to section 3.3. step 1,

$$
\begin{equation*}
\mu=\binom{0\left(-\frac{\dot{x}_{3}}{\left(1-\left(\dot{x}_{3}\right)^{2}\right)^{\frac{3}{2}}} d x_{2} \wedge d \dot{x}_{3}+\eta d \dot{x}_{3}\right) \wedge \frac{d}{d t}}{0} \tag{21}
\end{equation*}
$$

Step 2 yields $\eta=\frac{x_{2} \dot{x}_{3}}{\left(1-\dot{x}_{3}\right)^{\frac{3}{2}}}+\sigma\left(\dot{x}_{3}\right)$. For step 3 we set $M=\left(\begin{array}{cc}1 & m_{12} \frac{d}{d t} \\ 0 & 1\end{array}\right)$ which yields $m_{12}=-\left(\frac{x_{2}}{\sqrt{1-\left(\dot{x}_{3}\right)^{2}}}+\sigma_{1}\left(\dot{x}_{3}\right)\right)$ with $\sigma_{1}$ a primitive of $\sigma$. Thus, $d(M \omega)=0$ and setting $\left(d y_{1} d y_{2}\right)^{T}=M \omega$, one obtains

$$
\begin{equation*}
y_{1}=x_{1}-\frac{\dot{x}_{1}}{\dot{x}_{2}} x_{2}+\sigma_{2}\left(\dot{x}_{3}\right), \quad y_{2}=x_{3} \tag{22}
\end{equation*}
$$

where $\sigma_{2}\left(\dot{x}_{3}\right)$ is an arbitrary meromorphic function (a primitive of $\sigma_{1}$ ). By inversion of 22 we get

$$
\begin{align*}
& x_{1}=y_{1}-\arcsin \left(\dot{y}_{2}\right) \frac{\sqrt{1-\left(\dot{y}_{2}\right)^{2}}}{\ddot{y}_{2}}\left(\dot{y}_{1}-\sigma_{1}\left(\dot{y}_{2}\right) \ddot{y}_{2}\right)-\sigma_{2}\left(\dot{y}_{2}\right) \\
& x_{2}=-\frac{\sqrt{1-\left(\dot{y}_{2}\right)^{2}}}{\ddot{y}_{2}}\left(\dot{y}_{1}-\sigma_{1}\left(\dot{y}_{2}\right) \ddot{y}_{2}\right)  \tag{23}\\
& x_{3}=y_{2}
\end{align*}
$$

### 5.2 Academic Example: Euler-Lagrange Operator

We consider once more the example (18). We have

$$
\begin{equation*}
\frac{\partial F}{\partial \dot{x}}=\left(-\dot{x}_{2}^{-1} \cos \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right), \dot{x}_{1} \dot{x}_{2}^{-2} \cos \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right), 1\right), \mathcal{E}_{F}=\left(\eta_{1}, \eta_{2}, 0\right) \tag{24}
\end{equation*}
$$

with $\eta_{1}=-\frac{\ddot{x}_{2}}{\dot{x}_{2}^{2}} \cos \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right)-\frac{\ddot{x}_{1} \dot{x}_{2}-\dot{x}_{1} \ddot{x}_{2}}{\dot{x}_{2}^{3}} \sin \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right)$ and
$\eta_{2}=-\frac{\ddot{x}_{1} \dot{x}_{2}-2 \dot{x}_{1} \ddot{x}_{2}}{\dot{x}_{2}^{3}} \cos \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right)+\frac{\dot{x}_{1}\left(\ddot{x}_{1} \dot{x}_{2}-\dot{x}_{1} \ddot{x}_{2}\right)}{\dot{x}_{2}^{4}} \sin \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right)$.
The first two equations of (16), with $r_{1}=r_{2}=2$, read

$$
\begin{equation*}
-\frac{1}{\dot{x}_{2}} \cos \left(\frac{\dot{x}_{1}}{\dot{x}_{2}}\right)\left(\frac{\partial \varphi_{1}}{\partial \dot{y}_{j}}-\frac{\dot{x}_{1}}{\dot{x}_{2}} \frac{\partial \varphi_{2}}{\partial \dot{y}_{j}}\right)+\frac{\partial \varphi_{3}}{\partial \dot{y}_{j}}=0, \quad j=1,2 \tag{25}
\end{equation*}
$$

If we assume that $\frac{\partial \varphi_{3}}{\partial \ddot{y}_{j}}=\frac{\partial \varphi_{3}}{\partial \ddot{y}_{j}}=0, \quad j=1,2$ and introduce the variable

$$
\begin{equation*}
\psi=\frac{\dot{x}_{1}}{\dot{x}_{2}} \tag{26}
\end{equation*}
$$

with $\frac{\partial}{\partial \ddot{y}} \psi=0$ we obtain from 25

$$
\frac{\partial \varphi_{1}}{\partial \ddot{y}_{j}}-\psi \frac{\partial \varphi_{2}}{\partial \ddot{y}_{j}}=\frac{\partial}{\partial \ddot{y}_{j}}\left(\varphi_{1}-\psi \varphi_{2}\right)=0, \quad j=1,2
$$

Setting $\kappa(y, \dot{y})=\varphi_{1}-\psi \varphi_{2}$, we get

$$
\begin{equation*}
\dot{\kappa}=\dot{\varphi}_{1}-\psi \dot{\varphi}_{2}-\dot{\psi} \varphi_{2}=-\dot{\psi} \varphi_{2} \tag{27}
\end{equation*}
$$

Using the definition of $\kappa$ and 27 we obtain:

$$
\begin{equation*}
\varphi_{1}=\kappa-\frac{\dot{\kappa} \sqrt{1-\dot{\varphi}_{3}}}{\ddot{\varphi}_{3}} \arcsin \left(\dot{\varphi}_{3}\right), \quad \varphi_{2}=-\frac{\dot{\kappa}}{\ddot{\varphi}_{3}} \sqrt{1-\dot{\varphi}_{3}}, \quad \varphi_{3}=\varphi_{3}(y) \tag{28}
\end{equation*}
$$

Choosing $\varphi_{3}=y_{2}, \kappa=y_{1}$, we arrive at the invertible transformation

$$
x_{1}=\varphi_{1}=y_{1}-\frac{\dot{y}_{1}}{\ddot{y}_{2}} \sqrt{1-\dot{y}_{2}^{2}} \arcsin \left(\dot{y}_{2}\right), x_{2}=\varphi_{2}=-\frac{\dot{y}_{1}}{\ddot{y}_{2}} \sqrt{1-\dot{y}_{2}^{2}},
$$

with $x_{3}=\varphi_{3}=y_{2}$, which gives the same formula as 23 with $\sigma_{1}=$ $\sigma_{2}=0$. Hence $\left(y_{1}, y_{2}\right)$ is indeed a flat output, which implies that the remaining equations of (16) are satisfied.

### 5.3 An Example Proposed by P. Rouchon

Consider the implicit control system

$$
\begin{equation*}
F(x, \dot{x})=\dot{x}_{1} \dot{x}_{3}-\left(\dot{x}_{2}\right)^{2}=0 \tag{29}
\end{equation*}
$$

We thus have $\frac{\partial F}{\partial x}=\left(\begin{array}{lll}0 & 0 & 0\end{array}\right), \frac{\partial F}{\partial \dot{x}}=\left(\begin{array}{lll}\dot{x}_{3} & -2 \dot{x}_{2} & \dot{x}_{1}\end{array}\right)$ and

$$
\mathcal{E}_{F}=\frac{\partial F}{\partial x}-\frac{d}{d t}\left(\frac{\partial F}{\partial \dot{x}}\right)=-\frac{d}{d t}\left(\frac{\partial F}{\partial \dot{x}}\right)=\left(\begin{array}{lll}
-\ddot{x}_{3} & 2 \ddot{x}_{2} & -\ddot{x}_{1}
\end{array}\right) .
$$

The lowest possible choice of $\left(r_{1}, r_{2}\right)$ in Theorem 4 is $r_{1}=r_{2}=1$. However, there is no solution of (16) for these values, and we choose $r_{1}=r_{2}=2$. The two first equations of read

$$
\begin{equation*}
\dot{\varphi}_{3} \frac{\partial \varphi_{1}}{\partial \ddot{y}_{j}}-2 \dot{\varphi}_{2} \frac{\partial \varphi_{2}}{\partial \ddot{y}_{j}}+\dot{\varphi}_{1} \frac{\partial \varphi_{3}}{\partial \ddot{y}_{j}}=0, \quad j=1,2 \tag{30}
\end{equation*}
$$

We divide (30) by $\dot{\varphi}_{3}$ to obtain

$$
\begin{equation*}
\frac{\partial \varphi_{1}}{\partial \ddot{y}_{j}}-2 \psi \frac{\partial \varphi_{2}}{\partial \ddot{y}_{j}}+\psi^{2} \frac{\partial \varphi_{3}}{\partial \ddot{y}_{j}}=0, \quad j=1,2 \tag{31}
\end{equation*}
$$

where, taking account of the system equation 29,

$$
\begin{equation*}
\psi=\frac{\dot{\varphi}_{2}}{\dot{\varphi}_{3}}=\sqrt{\frac{\dot{\varphi}_{1}}{\dot{\varphi}_{3}}} \tag{32}
\end{equation*}
$$

If we assume that $\psi$ doesn't depend on $\ddot{y}_{1}$ and $\ddot{y}_{2}$, equation reads $\frac{\partial}{\partial \ddot{y}_{j}}\left(\varphi_{1}-2 \psi \varphi_{2}+\psi^{2} \varphi_{3}\right)=0$, for $j=1,2$. In other words, there exists a function $\kappa$ satisfying $\frac{\partial \kappa}{\partial \ddot{y}_{j}}=0$ for $j=1,2$, such that

$$
\begin{equation*}
\varphi_{1}-2 \psi \varphi_{2}+\psi^{2} \varphi_{3}=\kappa \tag{33}
\end{equation*}
$$

Differentiating the latter relation with respect to $t$, and taking into account the relation $\dot{\varphi}_{1}-2 \psi \dot{\varphi}_{2}+\psi^{2} \dot{\varphi}_{3}=0$ obtained from 29) and 32, we get

$$
\begin{equation*}
\varphi_{2}-\psi \varphi_{3}=-\frac{\dot{\kappa}}{2 \dot{\psi}} \tag{34}
\end{equation*}
$$

We again differentiate the latter relation with respect to $t$ to obtain

$$
\begin{equation*}
\varphi_{3}=\frac{\ddot{\kappa} \dot{\psi}-\dot{\kappa} \ddot{\psi}}{2 \dot{\psi}^{3}} \tag{35}
\end{equation*}
$$

thanks to $\dot{\varphi}_{2}-\psi \dot{\varphi}_{3}=0$ from (32). Thus, solving the system (33)-35), we immediately obtain

$$
\begin{align*}
& \varphi_{1}=\kappa-\psi \frac{\dot{\kappa}}{\dot{\psi}}+\psi^{2}\left(\frac{\ddot{\kappa} \dot{\psi}-\dot{\kappa} \ddot{\psi}}{2 \dot{\psi}^{3}}\right) \\
& \varphi_{2}=-\frac{\dot{\kappa}}{2 \dot{\psi}}+\psi\left(\frac{\ddot{\kappa} \dot{\psi}-\dot{\kappa} \psi}{2 \dot{\psi}^{3}}\right)  \tag{36}\\
& \varphi_{3}=\frac{\ddot{\kappa} \psi-\dot{\kappa} \ddot{\psi}}{2 \dot{\psi}^{3}}
\end{align*}
$$

where $\kappa$ and $\psi$ are arbitrary functions of $y_{1}, y_{2}, \dot{y}_{1}, \dot{y}_{2}$.
Note that choosing $\kappa=y_{1}$ and $\psi=y_{2}$ yields, after inversion of with (32):

$$
y_{1}=x_{1}-2 x_{2} \frac{\dot{x}_{2}}{\dot{x}_{3}}+x_{3} \frac{\dot{x}_{1}}{\dot{x}_{3}}, \quad y_{2}=\frac{\dot{x}_{2}}{\dot{x}_{3}},
$$

which is similar to the solution obtained by F. Ollivier ${ }^{5}$
Similarly, the solution of K. Schlacher and M. Schöberl [29] may be

[^2]recovered by posing $\kappa=y_{1}-y_{2} \frac{\dot{y}_{1}}{\dot{y}_{2}}$ and $\psi=\frac{\dot{y}_{1}}{2 \dot{y}_{2}}$ which, again after inversion of (36) with (32), yields:
$$
y_{1}=x_{1}-x_{3} \frac{\dot{x}_{1}}{\dot{x}_{3}}, \quad y_{2}=x_{2}-x_{3} \frac{\dot{x}_{2}}{\dot{x}_{3}} .
$$

## 6 Conclusion

In this survey we presented two dual approaches to flatness necessary and sufficient conditions, one based on the integration of 1 -forms and the second based on the integration of a set of PDEs involving a generalized Euler-Lagrange operator. Their complexity is compared on examples.

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[^0]:    ${ }^{3}$ From now on, $\bar{x} \bar{y}, \ldots$ stand for the sequences of jets of infinite order of $x, y, \ldots$

[^1]:    ${ }^{4}$ Note that if $\Phi$ is a meromorphic mapping from $\mathfrak{Y}$ to $\mathfrak{X}$, the (backward) image by $\Phi$ of a 1 -form is defined in the same way as in the finite dimensional context.

[^2]:    ${ }^{5}$ personal communication

