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Time-Optimal Control of a Particle in a Dielectrophoretic System

Dong Eui Chang, Nicolas Petit, and Pierre Rouchon

Abstract—We study the time-optimal control of a particle in a dielectrophoretic system. This system consists of a time-varying nonuniform electric field which acts upon the particle by creating a dipole within it. The interaction between the induced dipole and the electric field generates the motion of the particle. The control is the voltage on the electrodes which induces the electric field. Since we are considering the motion of a particle on an invariant line in a chamber filled with fluid flowing at low Reynolds number, the dynamics have a two dimensional state; one for the particle position and the other for the induced dipole moment. In regard to time-optimal control, we address the issue of existence and uniqueness of optimal trajectories, and explicitly compute the optimal control and the corresponding minimum time. Finally, we cast our analysis in the framework of symplectic reduction theory in order to provide geometric insight into the problem.

Index Terms—Biotechnology, dielectrophoresis, nanotechnology, time-optimal control.

I. INTRODUCTION

WE STUDY the time-optimal control of the following system:

$$\dot{x} = yu + \alpha u^2 \quad (1)$$

$$\dot{y} = -cy + u \quad (2)$$

with the state $(x, y) \in \mathbb{R}^2$ and the single control u satisfying

$$x(0) = x_0 = \text{given} \quad y(0) = 0 \quad (3)$$

$$x(t_f) = x_f = \text{given} \quad y(t_f) = \text{free} \quad (4)$$

$$|u| \leq 1 \quad (5)$$

where the parameters α and c satisfy

$$\alpha < 0 \quad c > 0. \quad (6)$$

These dynamics describe, after a nonlinear change of coordinates, the motion of a neutrally buoyant and neutrally charged particle on an invariant line in a chamber filled with fluid flowing at low Reynolds number and with a parallel electrode array at the bottom of the chamber. The existence of the invariant line is due to symmetry in the arrangement of electrodes and the boundary potential on electrodes. The motion is created by the interaction between a nonuniform electric field and the dipole moment

induced in the particle. This motion is called dielectrophoresis (DEP) [11]. Dielectrophoresis has wide applications in nano/bio-technology, in particular, in manipulating, separating and identifying nano/bio-particles [7], [8].

A brief explanation of the dynamics is in order. The variable x describes the displacement of the particle. The variable y describes the exponentially decaying part of the induced dipole moment. Voltage u is given on every other electrode and $(-u)$ on the others. Parameters α and c depend on the permittivities and conductivities of the particle and the fluid medium. The positivity of c is imposed by physics, but the negativity of α is arbitrary. As one is not interested in the final value of the induced dipole moment, the final value of y is free in (4).

Our goal in this paper is to study the minimum time trajectories of this system. A complete solution of this problem in a general setup—for a set of particles in a three-dimensional space, for instance—would allow significant improvement in DEP-based devices for particle analysis such as detecting cancers cells and separating different cells. However, here we address a simple case, which still includes the key feature of dielectrophoresis. Various control problems on dielectrophoresis in nano/biotechnology are suggested in [5].

This paper is organized as follows. We first overview the main results of the paper. Second, we derive the dynamics from physics. Third, we study the nonexistence of Lebesgue measurable time-optimal control for $x_f < x_0$ even though the target point is reachable. Fourth, we show that $(1 + \alpha c) > 0$ is the necessary and sufficient condition, under assumption (6), for the existence of time-optimal controls when $x_f > x_0$. Fifth, we address the issue of uniqueness of time-optimal control when $x_f > x_0$. We find a condition on α and c which guarantees the uniqueness of optimal trajectories, and compute the minimum time and the optimal control for a given target point x_f . For the case that the uniqueness condition is not satisfied, we give a constructive algorithm with which we can easily find a time-optimal control. Sixth, we make a discussion on the case where the optimal trajectories derived above are still valid in the presence of a state constraint on x such as $x \geq 0$ or $x \leq 0$. Seventh, we give geometric insight into the problem by putting the previous analysis in the picture of symplectic reduction theory. Finally, we perform some simulations to demonstrate the result.

II. OVERVIEW OF MAIN RESULTS

Time optimal trajectories satisfy the following dynamics:

$$\dot{x} = yu + \alpha u^2 \quad (7)$$

$$\dot{y} = -cy + u \quad (8)$$

$$\dot{\lambda} = c\lambda - u \quad (9)$$

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with conditions (3)–(6) and

$$\lambda(0) = \lambda_0 = \text{to be found} \quad \lambda(t_f) = 0 \quad (10)$$

where

$$u(t) = \arg \max_{|v| \leq 1} \tilde{H}(y(t), \lambda(t), v) \quad (11)$$

with

$$\tilde{H}(y, \lambda, v) = \alpha u^2 + (y + \lambda)u - cy\lambda. \quad (12)$$

The control u satisfying (11) is given by

$$u = \begin{cases} +1, & \text{if } y + \lambda > -2\alpha \\ \frac{y+\lambda}{-2\alpha}, & \text{if } |y + \lambda| \leq -2\alpha \\ -1, & \text{if } y + \lambda < 2\alpha \end{cases}$$

because $\alpha < 0$. A symplectic reduction picture is hidden in (7)–(12).

When $x_f < x_0$, there are no Lebesgue measurable time-optimal control functions resulting in $x(t_f) = x_f$ even though x_f is reachable.

We now consider the case of $x_f > x_0$. Time-optimal control law exists if and only if the parameters, α and c satisfy $(1 + \alpha c) > 0$, which is assumed in the following of this section. Define the open interval Λ by

$$\Lambda = (-2\alpha\sqrt{1 + \alpha c}, -2\alpha\sqrt{(1 + \alpha c)/(-\alpha c)}) \quad (13)$$

if $(1 + 2\alpha c) \leq 0$, and

$$\begin{aligned} \Lambda &= \Lambda_1 \cup \Lambda_2 \\ &= (-2\alpha\sqrt{1 + \alpha c}, -2\alpha) \cup [-2\alpha, 1/c) \end{aligned} \quad (14)$$

if $(1 + 2\alpha c) > 0$. For the sake of convenience, let us define four sentences as follows:

- P1 := $[(1 + 2\alpha c) > 0] \wedge [\lambda_0 \in \Lambda_1]$;
- P2 := $[(1 + 2\alpha c) > 0] \wedge [\lambda_0 \in \Lambda_2]$;
- Q := $[(1 + 2\alpha c) \leq 0] \wedge [\lambda_0 \in \Lambda]$;
- R := $[(3 + 4\alpha c) \geq 0] \wedge [\lambda_0 \in \Lambda]$;

where \wedge is the logical connective AND. Let us define a strictly increasing onto function $X_1 : \Lambda \rightarrow (0, \infty)$ by

$$X_1(\lambda_0) = \begin{cases} \text{equation (63),} & \text{if P1} \vee \text{Q} \\ \text{equation (67),} & \text{if P2} \end{cases} \quad (15)$$

where \vee is the logical connective OR. Let us define another function T_1 on Λ by

$$T_1(\lambda_0) = \begin{cases} \text{equation (73),} & \text{if P1} \vee \text{Q} \\ \text{equation (74),} & \text{if P2.} \end{cases} \quad (16)$$

We call a trajectory of (7)–(9) satisfying (3)–(5) and (10)–(11), an *extremal*. Let us call an arc of an extremal a *basic arc* if the projection of the arc onto the $y - \lambda$ plane starts from

$0 \times \Lambda$ (respectively, $0 \times (-\Lambda)$) and ends on $\Lambda \times 0$ (respectively, $(-\Lambda) \times 0$), going through the first (respectively, third) quadrant of the $y - \lambda$ plane. We call an extremal an *n-shot extremal* with $n \in \mathbb{N}$ if the maximum number of basic arcs in the extremal is n .

To decompose extremals into finite arcs when $[P1 \vee Q]$, let us introduce some notation. An arc associated with the linear control $u = (y + \lambda)/(-2\alpha)$ on a time interval of length $\Delta t_{AB}(\lambda_0)$ in (71), is denoted by $\gamma_L^{\lambda_0}$. Let $\gamma_+^{\lambda_0}$ (respectively, $\gamma_-^{\lambda_0}$) denote an arc with $u = 1$ (respectively, $u = -1$) on a time interval of length $\Delta t_{BC}(\lambda_0)$ in (72). Define two arcs $\Gamma_{\pm}^{\lambda_0}$ by the concatenation

$$\Gamma_{\pm}^{\lambda_0} = \gamma_L^{\lambda_0} \star \gamma_{\pm}^{\lambda_0} \star \gamma_L^{\lambda_0} \quad (17)$$

where the concatenation \star is defined such that the leftmost one comes first and the rightmost one comes last. An arc with the linear control $u = (y + \lambda)/(-2\alpha)$ is called an *idling arc* if its projection (y, λ) starts from the positive (respectively, negative) y -axis, goes through the fourth (respectively, second) quadrant in the $y - \lambda$ plane, and finally ends at the negative (respectively, positive) λ -axis. An idling arc is denoted by γ_{idling} ; see Fig. 8. Its duration T_{idling} is given in (76). Hence, when $[P1 \vee Q]$, we can express n -shot extremals as $\Gamma_{\pm}^{\lambda_0, n}$ for $n \in \mathbb{N}$, which are defined as follows:

$$\Gamma_{\pm}^{\lambda_0, 1} = \Gamma_{\pm}^{\lambda_0} \quad (18)$$

$$\Gamma_{\pm}^{\lambda_0, k} = \begin{cases} \Gamma_{\pm}^{\lambda_0, k-1} \star \gamma_{\text{idling}} \star \Gamma_{\mp}^{\lambda_0}, & \text{if } k \text{ is even} \\ \Gamma_{\pm}^{\lambda_0, k-1} \star \gamma_{\text{idling}} \star \Gamma_{\pm}^{\lambda_0}, & \text{if } k \text{ is odd} \end{cases} \quad (19)$$

for $k \geq 2$ with

$$\lambda(0) = \begin{cases} +\lambda_0, & \text{for } \Gamma_+^{\lambda_0, m} \\ -\lambda_0, & \text{for } \Gamma_-^{\lambda_0, m} \end{cases}$$

for $m \geq 1$.

We now discuss the existence and uniqueness of optimal trajectories. If $(3 + 4\alpha c) \geq 0$, there exist exactly two time-optimal trajectories for $x_f > x_0$, and they are basic arcs. Here is the procedure of constructing them.

[A.1.] Find $\lambda_0 = X_1^{-1}(x_f - x_0) \in \Lambda$.

[A.2.] Set $\lambda(0) = \pm\lambda_0$.

[A.3.] The minimum time cost is $T_1(\lambda_0)$ and the optimal trajectories γ_{opt} are

$$\gamma_{\text{opt}} = \begin{cases} \Gamma_{\pm}^{\lambda_0}, & \text{if R} \wedge \neg \text{P2} \\ \text{basic arc with } u = \pm 1, & \text{if P2} \end{cases}$$

where \neg is the logical connective NOT.

If $(3 + 4\alpha c) < 0$, we do not have any general proof of the uniqueness of optimal control. However, we have a finite procedure of finding all optimal control laws for $x_f > x_0$ as follows.

[B.1.] Define two sequences, for $k \in \mathbb{N}$

$$\begin{aligned} \lambda_{0, k} &= X_1^{-1} \left(\frac{x_f - x_0}{k} \right) \\ T_k &= k T_1(\lambda_{0, k}) + (k - 1) T_{\text{idling}}. \end{aligned}$$

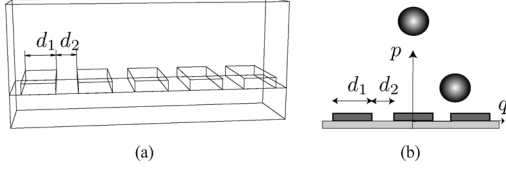


Fig. 1. Dielectrophoretic system is a chamber filled with a fluid medium where there is a parallel array of electrodes at the bottom.

[B.2.] Find $n = \arg \min_k \{T_k : k \in \mathbb{N}\}$ (such an n always exists and is less than $1 + (T_1/T_{\text{idling}})$).

[B.3.] Set $\lambda(0) = \pm \lambda_{0,n}$.

[B.4.] The minimum time cost is T_n and the corresponding optimal trajectories are $\Gamma_{\pm}^{\lambda_{0,n}}$.

If there are j integers in step [B.2.] to give the minimum time, then there are exactly $2j$ time-optimal trajectories. We remark that each basic arc in the k -shot extremal $\Gamma_{\pm}^{\lambda_{0,k},k}$ equally contributes $(x_f - x_0)/k$ to the increment of x and the idling arcs in between make no contribution.

In Section VI-H, we discuss the possibility that some optimal trajectories derived previously become optimal even with state constraints $x \geq 0$ or $x \leq 0$. In Section VII, we draw a symplectic reduction picture hidden in the problem and our solution.

III. DERIVATION OF DYNAMICS

We briefly derive the dynamics in (1) and (2), and explain the conditions in (3)–(6); see [4] for more details. Consider a neutrally charged particle in a chamber with a fluid medium and a parallel electrode array at the bottom as in Fig. 1(a) where d_1 is the width of each electrode, and d_2 is the width of the gap between two electrodes. As the electrodes are very long compared with the size of particles, we may assume that there are infinite number of infinitely long electrodes. Due to this symmetry, we can consider the motion of the particle in the vertical plane as in Fig. 1(b).

Let $(q, p) \in \mathbb{R}^2$ be the coordinates in Fig. 1(b). We give the boundary voltage

$$V_{\text{bd}}(t) = V_0 \cdot u(t), \quad |u(t)| \leq 1$$

on every other electrode and $(-V_{\text{bd}}(t))$ on the others. This creates potential $V(q, p, t)$ in $\{p \geq 0\}$. The electric field vector $\mathbf{E}(q, p, t) \in \mathbb{R}^2$ in $\{p \geq 0\}$ is given by

$$\mathbf{E}(q, p, t) = -\nabla V(q, p, t).$$

This electric field induces a dipole moment \mathbf{m} in a single-layered spherical particle as follows:

$$\mathbf{m}(q, p, t) = g(t) * \mathbf{E}(q, p, t)$$

where $*$ denotes the usual convolution operator with respect to time t and the Laplace transform $G(s)$ of the (transfer) function $g(t)$ is given by

$$G(s) = a + \frac{b}{s + c}$$

where

$$a = 4\pi r^3 \epsilon_m (\epsilon_p - \epsilon_m) / (\epsilon_p + 2\epsilon_m)$$

$$b = a \left(\frac{\sigma_p - \sigma_m}{\epsilon_p - \epsilon_m} - \frac{\sigma_p + 2\sigma_m}{\epsilon_p + 2\epsilon_m} \right) \quad (20)$$

$$c = (\sigma_p + 2\sigma_m) / (\epsilon_p + 2\epsilon_m) \quad (21)$$

where r is the radius of the particle, ϵ_p (respectively, ϵ_m) is the permittivity of the particle (resp., medium) and σ_p (respectively, σ_m) is the conductivity of the particle (respectively, medium). The interaction between the electric field and the induced dipole moment creates a force \mathbf{F}_{dep} . It is called *dielectrophoretic force* and is given by

$$\mathbf{F}_{\text{dep}}(q, p, t) = (\mathbf{m}(q, p, t) \cdot \nabla) \mathbf{E}(q, p, t).$$

We restrict our interest to the motion of a particle on the p -axis because it can practically represent the vertical motion of all particles in the whole chamber. One can check that the dielectrophoretic force on the p -axis is parallel to this axis due to the symmetry in the boundary voltage. This vertical dielectrophoretic force on the p -axis is denoted by $F_{\text{dep}}(p, t)$. It is of the form

$$F_{\text{dep}}(p, t) = F(p)u(t)(g * u)(t)$$

where

$$F(p) = \frac{-\pi^3 V_0^2 e^{-\frac{\pi p}{2d}} (1 - e^{-\frac{\pi p}{2d}})}{2d^3 K^2 \left(\cos \left(\frac{d_2 \pi}{4d} \right) \right) \left[1 - 2e^{-\frac{\pi p}{2d}} \cos \left(\frac{d_1 \pi}{2d} \right) + e^{-\frac{\pi p}{d}} \right]^2}$$

where K is the complete elliptic function of the first kind and $d = (d_1 + d_2)/2$; see [4] for the derivation of F . Notice that $F(p) \leq 0$ on $p \geq 0$, $\lim_{p \rightarrow \infty} F(p) = 0$, and $F(p) = 0$ only at $p = 0$.

Let us assume that *the particle is neutrally buoyant and the medium fluid flows at low Reynolds number*. Thus, the gravitational force and the buoyant force cancel and the inertial term $m\ddot{p}$ is trivial. The only forces on the particle are the drag and the DEP force. Hence, the motion of the particle on the p -axis can be described by

$$f\dot{p} + F(p)u(t)(g * u)(t) = 0. \quad (22)$$

where $f > 0$ is the drag constant.

We assume that b in (20) is nonzero, which generically holds. Then, (1) and (2) come from (22) where x and y are defined by

$$x = \int_{\epsilon}^p \frac{-f}{bF(z)} dz$$

$$Y(s) = \frac{1}{s + c} U(s) \quad (23)$$

for $p \geq \epsilon$ where ϵ is a positive number and $Y(s)$ and $U(s)$ are the Laplace transforms of $y(t)$ and $u(t)$, and α is defined by

$$\alpha = a/b.$$

If a particle is close to the electrode, then additional physical/chemical forces other than the DEP force start to appear in the dynamics [7], [8], [11], so the parameter ϵ in (23) defines the region where the dynamics (22) is valid. Physically, y is the exponentially induced part of the dipole moment, so we have the initial condition $y(0) = 0$. As we are not interested in the final state of the induced dipole moment, we have $y(t_f) = \text{free}$.

Depending on the sign of b , the original region $\{p \geq \epsilon\}$ is mapped to $\{x \geq 0\}$ or $\{x \leq 0\}$. In this paper, we ignore this state constraint on x , allowing for x to be on the whole real line. In Section VI-H, we discuss the possibility that the time optimal trajectories without the state constraints remain optimal with the state constraints. In a future publication, we will address the optimization problem with the state constraint on x .

We also make the following assumption on the signs of parameters α and c

$$\alpha < 0 \quad c > 0.$$

The assumption $c > 0$ is imposed by physics; see (21). However, the condition $\alpha < 0$ is chosen for convenience. The case where $\alpha \geq 0$ is left for future work.

IV. NONEXISTENCE OF OPTIMAL CONTROL FOR $x_f < x_0$

We will show that there are no time-optimal (Lebesgue) measurable control functions for $x_f < x_0$ even though x_f is reachable.

Fix a $T > 0$. Let us define a sequence of functions $\{u_n^T : [0, T] \rightarrow \pm 1\}_{n \in \mathbb{N}}$ as follows:

$$u_n^T(t) = \text{sign}(\sin(2\pi nt/T)). \quad (24)$$

It is straightforward to prove the following lemma.

Lemma 4.1: For any continuous function f on $[0, T]$

$$\lim_{n \rightarrow \infty} \int_0^t f \cdot u_n^T = 0 \quad (25)$$

uniformly in $t \in [0, T]$.

The substitution of u in (2) to (1) yields

$$\dot{x} = y\dot{y} + cy^2 + \alpha u^2.$$

It follows that

$$\begin{aligned} x(T) - x(0) &= \int_0^T \dot{x} dt \\ &= \frac{1}{2}y(T)^2 + c \int_0^T y(\tau)^2 d\tau + \alpha \int_0^T u(\tau)^2 d\tau \end{aligned} \quad (26)$$

$$\begin{aligned} &\geq \alpha \int_0^T u(\tau)^2 d\tau \\ &\geq \alpha T \end{aligned} \quad (27)$$

because $\alpha < 0$, $x(0) = 0$, and $|u| \leq 1$. This implies that $x(0) + \alpha T$ is a lower bound of $x(T)$ for any admissible control u . Let $(x_n(t), y_n(t)), 0 \leq t \leq T$ be the solution to (1) and (2) with control u_n^T in (24). In particular

$$y_n(t) = e^{-ct} \int_0^t e^{c\tau} u_n^T(\tau) d\tau. \quad (28)$$

Given $\epsilon > 0$, by Lemma 4.1, there exists $N \in \mathbb{N}$ such that

$$|y_n(t)| \leq \sqrt{\epsilon} \quad (29)$$

for all $t \in [0, T]$ and all $n \geq N$. By (26), (29) and the definition of u_n^T , we have

$$|x_n(T) - x(0) - \alpha T| \leq \left(\frac{1}{2} + cT\right) \epsilon$$

for all $n > N$. Hence

$$\lim_{n \rightarrow \infty} x_n(T) = x(0) + \alpha T.$$

We have constructed a sequence $\{u_n^T\}$ of control laws such that the corresponding $\{x_n(T)\}$ converges to the lower bound $x(0) + \alpha T$ of the reachable point of x in time T . Notice that the sequence of functions $\{u_n^T\}$ does not converge to a measurable function. The following lemma addresses this issue.

Lemma 4.4: For a given $T > 0$, there exists no measurable control function $u : [0, T] \rightarrow [-1, 1]$ such that $x(T) = x(0) + \alpha T$.

Proof: The proof is by contradiction. Suppose that there is a measurable function $u : [0, T] \rightarrow [-1, 1]$ such that $x(T) = x(0) + \alpha T$. By (27), it follows that

$$y(t) = 0 \quad \text{a.e. on } [0, T] \quad (30)$$

$$u(t) = \pm 1 \quad \text{a.e. on } [0, T]. \quad (31)$$

Hence, for almost all $t \in [0, T]$

$$\begin{aligned} 0 = y(t) &= \int_0^t \dot{y}(s) ds \\ &= \int_0^t (-cy(s) + u(s)) ds = \int_0^t u(s) ds. \end{aligned}$$

By [6, Th. 4.9], the function $t \mapsto \int_0^t u(s) ds$ is continuous. It follows that

$$\int_0^t u(s) ds = 0$$

for all $t \in [0, T]$. By [6, 4.11], $u(t) = 0$ for almost all $t \in [0, T]$ which contradicts (31). Thus, there exists no such measurable function u that produces $x(T) = x(0) + \alpha T$. ■

We have shown the following.

Claim 4.3: For $x_f < x_0$, the infimum of the time cost is $T = (x_0 - x_f)/(-\alpha)$, but there are no time-optimal (Lebesgue) measurable controls to reach x_f in time T .

Remark 4.4: It will become an interesting project to extend the technology in [10] and [3] in order to show the nonexistence of optimal trajectories in an alternative way.

Remark 4.5: For $x_f = x_0$, the control $u = 0$ with $t_f = 0$ is trivially the time-optimal control.

V. PONTRYAGIN MAXIMUM PRINCIPLE

We derive, from the Pontryagin Maximum Principle (PMP), a necessary condition which time-optimal trajectories must satisfy. Let us define the PMP Hamiltonian H for the time-optimal control as follows (see [12] and [2]):

$$H(x, y, \lambda_x, \lambda_y, u) = \lambda_x \alpha u^2 + (\lambda_x y + \lambda_y)u - cy\lambda_y \quad (32)$$

where (λ_x, λ_y) is a covector. Let

$$M^\circ(x, y, \lambda_x, \lambda_y) = \max_{|u| \leq 1} H(x, y, \lambda_x, \lambda_y, u). \quad (33)$$

The application of the PMP in [12] gives:

Theorem 5.1: Consider system (1), (2) with conditions (3)–(6). Let $u(t)$ be a time-optimal control and $(x(t), y(t))$ be the corresponding trajectory. Then, it is necessary that there exists a continuous covector $(\lambda_x(t), \lambda_y(t))$, which is not identically zero, such that

$$\dot{x} = yu + \alpha u^2 \quad (34)$$

$$\dot{y} = -cy + u \quad (35)$$

$$\dot{\lambda}_x = 0 \quad (36)$$

$$\dot{\lambda}_y = c\lambda_y - \lambda_x u. \quad (37)$$

Additionally, the following must be satisfied:

$$1. \quad u(t) = \arg \max_{|v| \leq 1} H(x(t), y(t), \lambda_x(t), \lambda_y(t), v) \quad (38)$$

$$2. \quad M^\circ(x(t), y(t), \lambda_x(t), \lambda_y(t)) = \text{constant} \geq 0 \quad (39)$$

$$3. \quad \lambda_y(t_f) = 0 \text{ (transversality condition).} \quad (40)$$

The boundary conditions of (x, y) and the signs of parameters are given in (3)–(6). The transversality condition in (40) comes from the free final boundary condition on y in (4).

Remark 5.2: Since $x(t) \in \mathbb{R}$, one can equivalently formulate the previous time-optimal control problem as follows:

$$\text{maximize } x(T)$$

for $T \geq 0$.

VI. ANALYSIS OF EXTREMALS

We study the dynamics in (34)–(37). We call the trajectories satisfying the dynamics and all the conditions in Theorem 5.1, extremals.

A. The Necessary Positivity of λ_x

By (36), $\lambda_x = \lambda_x(t)$ is constant in t , so there can be the following three cases:

$$\lambda_x = 0 \quad \lambda_x < 0 \quad \lambda_x > 0.$$

We will show that extremals exist only if $\lambda_x > 0$.

First, we assume $\lambda_x = 0$. Then (36) and (37) become a linear ordinary differential equation in (λ_x, λ_y) with $(\lambda_x(t_f), \lambda_y(t_f)) = (0, 0)$. By the uniqueness theorem of solutions of ordinary differential equations [1] or by direct computation, we have $(\lambda_x(t), \lambda_y(t)) = 0$ for all t . Hence, by the PMP, there are no optimal trajectories when $\lambda_x = 0$.

We now assume $\lambda_x < 0$. Let

$$\begin{aligned} \lambda &= \lambda_y / \lambda_x \\ \tilde{H}(y, \lambda, u) &= H(x, y, \lambda_x, \lambda_y, u) / \lambda_x \\ &= \alpha u^2 + (y + \lambda)u - cy\lambda. \end{aligned} \quad (41)$$

Then, (38) implies

$$u = \arg \min_{|v| \leq 1} \tilde{H}(y, \lambda, v).$$

One can compute (recall that $\alpha < 0$)

$$u(t) = -\text{sign}(y(t) + \lambda(t)).$$

With this control, the dynamics of (y, λ) can be written as

$$\begin{bmatrix} \dot{y} \\ \dot{\lambda} \end{bmatrix} = \begin{bmatrix} -c & 0 \\ 0 & c \end{bmatrix} \begin{bmatrix} y \\ \lambda \end{bmatrix} + \begin{bmatrix} -1 \\ 1 \end{bmatrix} \cdot \text{sign}(y + \lambda). \quad (42)$$

Suppose $\lambda(0) > 0$. Since $y(0) = 0$, $y(0) + \lambda(0) > 0$. One can compute the flow of (42)

$$y(t) = \frac{1}{c}(e^{-ct} - 1) \quad \lambda(t) = \lambda(0)e^{ct} + \frac{1}{c}(e^{ct} - 1).$$

It follows

$$y(t) + \lambda(t) = \lambda(0)e^{ct} + \frac{2}{c}(\cosh(ct) - 1) > 0.$$

Hence, $u(t) = -\text{sign}(y(t) + \lambda(t)) = -1$ for all $t \geq 0$, which implies that $\lambda(t) > 0$ for all $t \geq 0$. Hence, the transversality condition (40) cannot be satisfied. One can show in the similar manner that the transversality condition cannot be satisfied

when $\lambda(0) < 0$ or $\lambda(0) = 0$. Therefore, there are no optimal trajectories if $\lambda_x < 0$. Hence, we have proved the following claim.

Claim 6.1: Along every extremal, we have $\lambda_x > 0$.

B. A Necessary Condition on Parameters for the Existence of Optimal Trajectories

We derive a necessary condition on the parameters α and c so that optimal trajectories exist. By Claim 6.1, we assume $\lambda_x > 0$ in the following.

Let

$$\begin{aligned} \lambda &= \lambda_y / \lambda_x \\ \tilde{H}(y, \lambda, u) &= H(x, y, \lambda_x, \lambda_y, u) / \lambda_x \end{aligned} \quad (43)$$

$$\begin{aligned} &= \alpha u^2 + (y + \lambda)u - cy\lambda \\ M(y, \lambda) &= \max_{|v| \leq 1} \tilde{H}(y, \lambda, v) \end{aligned} \quad (44)$$

and

$$\begin{aligned} R_n &= \{(y, \lambda) \mid y + \lambda < 2\alpha\} \\ R_p &= \{(y, \lambda) \mid y + \lambda > -2\alpha\} \\ R_l &= \{(y, \lambda) \mid |y + \lambda| \leq -2\alpha\}. \end{aligned}$$

By (38) and (44)

$$u = \arg \max_{|v| \leq 1} \tilde{H}(y, \lambda, v).$$

Let us compute u in each region of R_n , R_p , and R_l and study the dynamics in each region.

First, we consider the case where $(y, \lambda) \in R_n$. Then, $u = -1$. The (y, λ) -dynamics become

$$\dot{y} = -cy - 1 \quad \dot{\lambda} = c\lambda + 1 \quad (45)$$

where the equilibrium at $(y, \lambda) = (-1/c, -1/c)$ is a saddle.

Second, if $(y, \lambda) \in R_p$, then $u = 1$. The (y, λ) -dynamics become

$$\dot{y} = -cy + 1 \quad \dot{\lambda} = c\lambda - 1 \quad (46)$$

where the equilibrium at $(y, \lambda) = (1/c, 1/c)$ is a saddle.

Finally, if $(y, \lambda) \in R_l$, then $u = (y + \lambda)/(-2\alpha)$. The Hamiltonian M becomes

$$M(y, \lambda) = \frac{(y + \lambda)^2}{-4\alpha} - cy\lambda \quad (47)$$

and the (x, y, λ) -dynamics become

$$\dot{x} = \frac{y^2 - \lambda^2}{-4\alpha} \quad (48)$$

$$\begin{bmatrix} \dot{y} \\ \dot{\lambda} \end{bmatrix} = A \begin{bmatrix} y \\ \lambda \end{bmatrix} \quad (49)$$

where

$$A = \begin{bmatrix} -(c + \frac{1}{2\alpha}) & -\frac{1}{2\alpha} \\ \frac{1}{2\alpha} & (c + \frac{1}{2\alpha}) \end{bmatrix}. \quad (50)$$

The matrix A satisfies

$$\text{tr } A = 0 \quad \det A = c(1 + \alpha c)/(-\alpha).$$

The type of the equilibrium at $(y, \lambda) = (0, 0)$ depends on the sign of $\det A$.

We now make qualitative phase portraits of the (y, λ) -dynamics in the following three different cases:

$$(1 + \alpha c) < 0 \quad (1 + \alpha c) = 0 \quad (1 + \alpha c) > 0.$$

Suppose $(1 + \alpha c) < 0$. Then, $\det A < 0$. The origin $(y, \lambda) = (0, 0)$ is a saddle point of (49). The (real) eigenvalues of A are given by

$$\mu_s(A) = \frac{\sqrt{\alpha c(1 + \alpha c)}}{\alpha} \quad \mu_u(A) = \frac{\sqrt{\alpha c(1 + \alpha c)}}{-\alpha}$$

and the corresponding eigenvectors are given by

$$\begin{aligned} v_s(A) &= (1, (\sqrt{-\alpha c} - \sqrt{-(\alpha c + 1)})^2)^T \\ v_u(A) &= (1, (\sqrt{-\alpha c} + \sqrt{-(\alpha c + 1)})^2)^T. \end{aligned}$$

Notice that $\mu_s(A) < 0$ and $\mu_u(A) > 0$. The second (or λ -) components of $v_s(A)$ and $v_u(A)$ satisfy

$$0 < v_{s,\lambda}(A) < 1 < v_{u,\lambda}(A) \quad (51)$$

because $(1 + \alpha c) < 0$. In addition, the two saddle points $\pm(1/c, 1/c)$ of the dynamics in (45) and (46) do not belong to $R_n \cup R_p$ but to R_l since $(1 + \alpha c) < 0$. Gathering the information in each region of R_n , R_p , and R_l , we can draw a phase portrait of the (y, λ) -dynamics, qualitatively, in Fig. 2. By (51), the slope of the unstable (respectively, stable) manifold of the origin is greater (respectively, smaller) than 1. As $y(0) = 0$ and $\lambda(t_f) = 0$, time-optimal trajectories must start from the λ -axis and ends on the y -axis in the y - λ plane. However, there are no such orbits in Fig. 2 because the stable and unstable manifolds of the saddle points $\pm(1/c, 1/c)$ prevent it in $R_n \cup R_p$. This implies that there are no optimal trajectories when $\lambda_x > 0$ and $(1 + \alpha c) < 0$.

If $(1 + \alpha c) = 0$, then the matrix A in (50) becomes

$$A = \begin{bmatrix} -\frac{c}{2} & \frac{c}{2} \\ -\frac{c}{2} & \frac{c}{2} \end{bmatrix}.$$

The integration of (49) gives $\lambda(t) = y(t) + c$ where c is a constant. In particular, the line $\lambda = y$ in R_l is a set of equilibria. The phase portrait in the y - λ plane is given in Fig. 3. One can see that no trajectories starting from the λ -axis reach the y -axis.

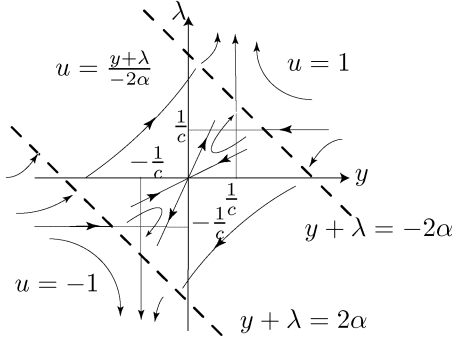


Fig. 2. Phase portrait in the y - λ plane when $\lambda_x > 0$ and $(1 + \alpha c) < 0$.

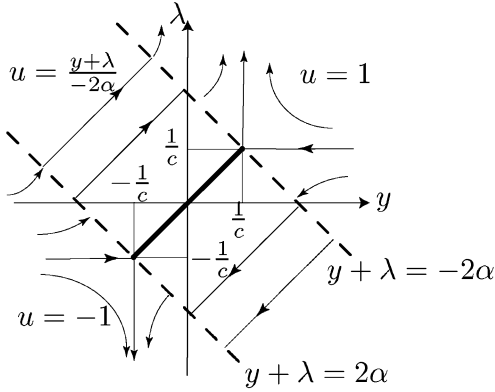


Fig. 3. Phase portrait in the y - λ plane when $\lambda_x > 0$ and $(1 + \alpha c) = 0$. The segment between $(-1/c, -1/c)$ and $(1/c, 1/c)$ is the set of equilibria.

Hence, there are no time-optimal trajectories when $\lambda_x > 0$ and $(1 + \alpha c) = 0$.

We have, so far, proved the following.

Claim 6.2: Time-optimal trajectories exist only if $(1 + \alpha c) > 0$.

When $(1 + \alpha c) > 0$, the fixed point $(y, \lambda) = (0, 0)$ of the dynamics (49) is a center as $\det A > 0$ with A in (50). In this case, there are two possible qualitatively different phase portraits of the (y, λ) -dynamics depending on the position of the λ -intercept of the switching line $y + \lambda = -2\alpha$ relative to the point $(0, 1/c)$ on the λ -axis. They are given in Fig. 4 depending on the sign of $(1 + 2\alpha c)$, with control

$$u = \begin{cases} +1, & \text{if } y + \lambda > -2\alpha \\ \frac{y+\lambda}{-2\alpha}, & \text{if } |y + \lambda| \leq -2\alpha \\ -1, & \text{if } y + \lambda < -2\alpha. \end{cases} \quad (52)$$

We will show that the (y, λ) -projection of optimal trajectories must be contained in the shaded region in Fig. 4. In the rest of this paper, we assume that $(1 + \alpha c) > 0$.

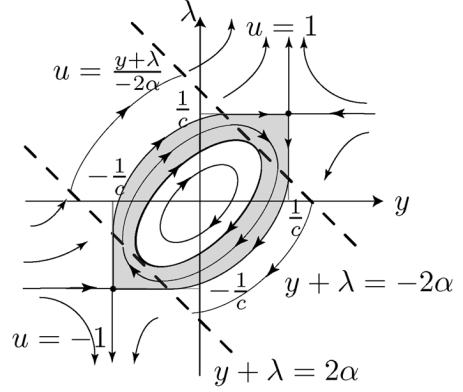
C. Discrete Symmetry

We will find $\mathbb{Z}_2 \times \mathbb{Z}_2$ symmetry in the dynamics. Define the following maps:

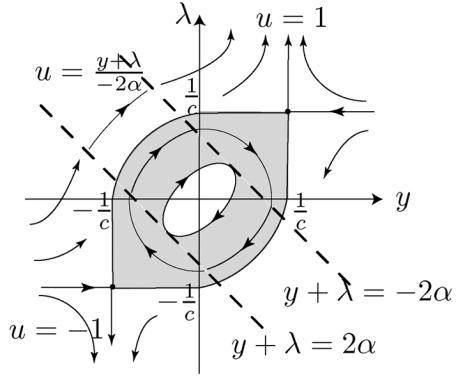
$$S_1 : (x, y, \lambda) \mapsto (x, \lambda, y) \quad (53)$$

$$S_2 : (x, y, \lambda) \mapsto (x, -\lambda, -y). \quad (54)$$

$$S_3 := S_1 \circ S_2 : (x, y, \lambda) \mapsto (x, -y, -\lambda). \quad (55)$$



(a)



(b)

Fig. 4. Phase portrait in the y - λ plane when $\lambda_x > 0$ and $(1 + \alpha c) > 0$. Depending on the sign of $(1 + 2\alpha c)$, the λ -intercept, -2α , of the switching line $y + \lambda = -2\alpha$ is greater or less than $1/c$. (a) $(1 + 2\alpha c) \leq 0$. (b) $(1 + 2\alpha c) > 0$.

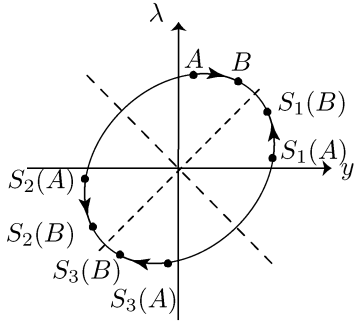
Denoting by $Z_l(x, y, \lambda)$ the vector field in (7)–(9) with the linear control $u = (y + \lambda)/(-2\alpha)$ on $\mathbb{R} \times R_l$, i.e., that in (48) and (49), we obtain

$$S_i \circ Z_l = -Z_l \circ S_i, \quad i = 1, 2. \quad (56)$$

The linear vector field Z_l is invariant under the reflections, S_1 and S_2 , up to the time-reversal, and it is invariant under the reflection S_3 without time-reversal. The region $\mathbb{R} \times R_l$ is invariant under Z_l , $i = 1, 2$. Notice this symmetry in the phase portraits in region R_l of Fig. 4. This symmetry gives us useful information as follows. Consider a trajectory of Z_l whose (y, λ) -projected image is contained in R_l as in Fig. 5. The duration Δt_{AB} from A to B along the trajectory with $u = (y + \lambda)/(-2\alpha)$ in the y - λ plane is the same as $\Delta t_{S_i(B)S_i(A)}$, $i = 1, 2$. Also, the corresponding (positive or negative) increments in x satisfy

$$\Delta x_{AB} = -\Delta x_{S_i(B)S_i(A)}, \quad i = 1, 2.$$

This implies that there cannot be any optimal trajectories in the white region surrounded by the shaded region in Fig. 4 because an arbitrary trajectory starting from the λ -axis and ending at the y -axis can be decomposed into parts, each of which is invariant under S_1 or S_2 , so $\Delta x = 0$ along the whole trajectory. However, the time-optimal control for $\Delta x = 0$ is $u = 0$ with


 Fig. 5. $\mathbb{Z}_2 \times \mathbb{Z}_2$ -symmetry in the linear region.

$t_f = 0$; see Remark 4.5. Therefore, all time-optimal trajectories are contained in the shaded regions in Fig. 4 because only the trajectories in the shaded region can satisfy $y(0) = 0$ and $\lambda(t_f) = 0$.

Let us consider the symmetry S_3 in (55). Let Z be the vector field in (7)–(9) on the whole domain, $\mathbb{R}^3 = \mathbb{R} \times (R_n \cup R_p \cup R_l)$. One can check

$$\begin{aligned} S_3(\mathbb{R} \times R_n) &= \mathbb{R} \times R_p \\ S_3(\mathbb{R} \times R_l) &= \mathbb{R} \times R_l \\ S_3 \circ Z &= Z \circ S_3 \\ \Delta t_{AB} &= \Delta t_{S_3(A)S_3(B)}. \end{aligned}$$

Hence, for example

$$\Delta x_{AB} = \Delta x_{S_3(A)S_3(B)}$$

in Fig. 5.

D. Definition, Monotonicity, and Positivity of X_1

Let us first define the basic region in the $y - \lambda$ plane. In each phase portrait in Fig. 4, we denote by \mathcal{A} the interior of the shaded region. The *basic region* \mathcal{B} is defined by

$$\mathcal{B} = \{(y, \lambda) \in \mathcal{A} \subset \mathbb{R}^2 \mid y \geq 0, \lambda \geq 0\}.$$

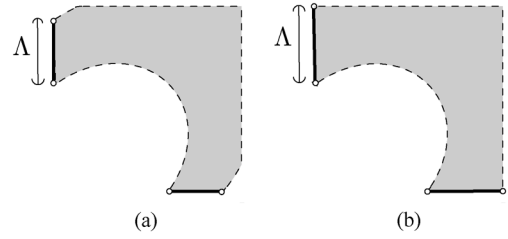
See Fig. 6 for an illustration of the basic region.

Given an extremal $\gamma(t) = (x(t), y(t), \lambda(t))$, $0 \leq t \leq T$, the arc $\gamma([t_1, t_2])$ with $0 \leq t_1 < t_2 \leq T$ is called a *basic arc* of γ if $y(t_1) = 0$, $\lambda(t_2) = 0$ and

$$\tilde{\pi}(\gamma([t_1, t_2])) \subset \mathcal{B} \cup \tilde{S}_3(\mathcal{B})$$

where $\tilde{\pi}(x, y, \lambda) = (y, \lambda)$ and $\tilde{S}_3(y, \lambda) = (-y, -\lambda)$.

We denote by Λ the open interval which is the intersection of the positive λ -axis with the basic region; see Fig. 6. It is given in (13) and (14), whose derivation will be made later. Let us construct a function X_1 of $\lambda_0 \in \Lambda$, which measures the (signed) increment of x along a basic arc starting with $\lambda(0) = \lambda_0$. The


 Fig. 6. Basic region and the domain Λ of the map X_1 . (a) $(1 + 2\alpha c) \leq 0$. (b) $(1 + 2\alpha c) > 0$.

construction of X_1 and Λ is made in the following two separate cases:

$$(1 + 2\alpha c) \leq 0 \quad (1 + 2\alpha c) > 0.$$

Case $(1 + 2\alpha c) \leq 0$: The domain Λ of X_1 is computed as follows:

$$\begin{aligned} \Lambda &= \{\lambda \in \mathbb{R}_+ \mid M(-\alpha, -\alpha) < M(0, \lambda) \\ &\quad < M(-(1 + 2\alpha c)/c, 1/c)\} \\ &= (-2\alpha\sqrt{1 + \alpha c}, -2\alpha\sqrt{(1 + \alpha c)/(-\alpha c)}) \end{aligned} \quad (57)$$

which is exactly (13), where the Hamiltonian M is given in (47). The switching line $\{y + \lambda = -2\alpha\}$ is tangent to the level set $\{M(y, \lambda) = M(-\alpha, -\alpha)\}$ at $(-\alpha, -\alpha)$. The level set $\{M(y, \lambda) = M(-(1 + 2\alpha c)/c, 1/c)\}$ goes through the intersection of the switching line $\{y + \lambda = -2\alpha\}$ and a stable manifold $\{\lambda = 1/c\}$ of $(1/c, 1/c)$; see Fig. 4(a). Notice in Fig. 4(a) that all the basic arcs are of form $\Gamma_{\pm}^{\lambda_0}$ defined in (17). Namely, basic arcs in the basic region are like arc $ABCD$ in Fig. 7(a). Let $A = (0, \lambda_0)$ in Fig. 7(a). Then, $D = (\lambda_0, 0)$ by the S_1 symmetry. The two points B and C in Fig. 7(a) are the intersection of the level set of M

$$D_1 = \{(y, \lambda) \mid M(y, \lambda) = M(0, \lambda_0)\} \quad (58)$$

and the switching line

$$D_2 = \{(y, \lambda) \mid y + \lambda = -2\alpha\}. \quad (59)$$

These points are symmetric to each other with respect to S_1 in (53) because D_1 and D_2 are invariant under S_1 . The y -coordinates of B and C are given by

$$y_B(\lambda_0) = -\alpha - \sqrt{\alpha^2 + (\lambda_0^2 - 4\alpha^2)/(-4\alpha c)} \quad (60)$$

$$y_C(\lambda_0) = -y_B(\lambda_0) - 2\alpha. \quad (61)$$

Due to the S_1 symmetry, the increments in x along AB and CD cancel each other, i.e.,

$$\Delta x_{AB} + \Delta x_{CD} = 0. \quad (62)$$

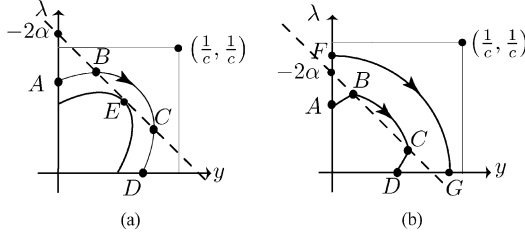


Fig. 7. Construction of the x -increment map X_1 . (a) $(1 + 2\alpha c) \leq 0$. (b) $(1 + 2\alpha c) > 0$.

Hence

$$\begin{aligned} X_1(\lambda_0) &= \Delta x_{BC} = \int_{y_B(\lambda_0)}^{y_C(\lambda_0)} \frac{\dot{x}}{\dot{y}} dy \\ &= \int_{y_B}^{y_C} \frac{y + \alpha}{-cy + 1} dy \\ &= \frac{1}{c} \left[-2\alpha - 2y_C(\lambda_0) \right. \\ &\quad \left. + \left(\alpha + \frac{1}{c} \right) \ln \left(\frac{1 + 2\alpha c + cy_C(\lambda_0)}{1 - cy_C(\lambda_0)} \right) \right] \end{aligned} \quad (63)$$

with $y_C(\lambda_0)$ in (61). Regarding X_1 as a function of y_C , one can see

$$\frac{dX_1}{dy_C} = \frac{2(y_C + \alpha)^2}{c \left(\frac{1}{c} + 2\alpha + y_C \right) \left(\frac{1}{c} - y_C \right)} \geq 0 \quad (64)$$

on $y_C(\Lambda) = (-\alpha, 1/c)$ because $(1 + \alpha c) > 0$. Since y_C in (61) is a strictly increasing function of λ_0 , it follows that X_1 is a strictly increasing function on Λ . Note that

$$\begin{aligned} \lim_{\lambda_0 \rightarrow -2\alpha\sqrt{1+\alpha c}} X_1(\lambda_0) &= \lim_{y_C \rightarrow -\alpha} X_1(y_C) = 0 \\ \lim_{\lambda_0 \rightarrow -2\alpha\sqrt{\frac{1+\alpha c}{-\alpha c}}} X_1(\lambda_0) &= \lim_{y_C \rightarrow 1/c} X_1(y_C) = +\infty. \end{aligned} \quad (65)$$

It follows that $X_1(\Lambda) = (0, \infty)$.

Case $(1 + 2\alpha c) > 0$: The domain Λ of X_1 is given by

$$\Lambda = \Lambda_1 \cup \Lambda_2 = (-2\alpha\sqrt{1+\alpha c}, -2\alpha) \cup [-2\alpha, 1/c] \quad (66)$$

where the left end $-2\alpha\sqrt{1+\alpha c}$ of Λ is computed from $M(0, \lambda) = M(-\alpha, -\alpha)$ and the right end $1/c$ is just the intersection of the stable manifold of the fixed point $(1/c, 1/c)$, with the λ -axis in Fig. 4(b). As $-2\alpha\sqrt{1+\alpha c} < -2\alpha < 1/c$, the decomposition of Λ into Λ_1 and Λ_2 is valid. Notice in Fig. 4(b) that there are two different kinds of basic arcs in the basic region. If $\lambda_0 \in \Lambda_1$, then the trajectory is like $ABCD$ in Fig. 7(b), i.e., $\Gamma_+^{\lambda_0}$ in (17). In this case, the control is given by $u = (y + \lambda)/(-2\alpha)$ on AB and CD and $u = +1$ on BC . If $\lambda_0 \in \Lambda_2$, then the trajectory is like FGD where the control is given by $u = +1$.

First, we consider the case where $\lambda_0 \in \Lambda_1$. Let y_B and y_C be the y -coordinates of the intersections points of D_1 and D_2 in

(58) and (59). The formulas of y_B and y_C are given in (60) and (61). Due to the S_1 symmetry, relation (62) holds. Hence

$$\begin{aligned} X_1(\lambda_0) &= \frac{1}{c} \left[-2\alpha - 2y_C(\lambda_0) \right. \\ &\quad \left. + \left(\alpha + \frac{1}{c} \right) \ln \left(\frac{1 + 2\alpha c + cy_C(\lambda_0)}{1 - cy_C(\lambda_0)} \right) \right] \end{aligned}$$

which coincides with (63). One can check that (64) is still valid on $y_C(\Lambda_1)$, and that y_C is a strictly increasing function of λ_0 on Λ_1 . Hence, X_1 is a strictly increasing function of λ_0 on Λ_1 .

We now consider the case that $\lambda_0 \in \Lambda_2$. We have

$$\begin{aligned} X_1(\lambda_0) &= \int_{y_F}^{y_G} \frac{\dot{x}}{\dot{y}} dy = \int_0^{\lambda_0} \frac{y + \alpha}{-cy + 1} dy \\ &= \frac{1}{c} \left[-\lambda_0 + \left(\alpha + \frac{1}{c} \right) \ln \left(\frac{1}{1 - c\lambda_0} \right) \right]. \end{aligned} \quad (67)$$

Notice the continuity of X_1 at $\lambda_0 = -2\alpha$ from (63) and (67). Since

$$\frac{dX_1}{d\lambda_0} = \frac{\lambda_0 + \alpha}{1 - c\lambda_0} > 0$$

for $\lambda_0 \in \Lambda_2$, the function $X_1(\lambda_0)$ is strictly increasing on Λ_2 . Notice that

$$\lim_{\lambda_0 \rightarrow 1/c} X_1(\lambda_0) = +\infty. \quad (68)$$

From (65) and (68), it follows that $X_1(\Lambda) = (0, \infty)$. We conclude that the map $X_1(\lambda_0)$ is strictly increasing on Λ and its image is $(0, \infty)$.

We have proved the following.

Claim 6.3: Irrespective of the sign of $(1 + 2\alpha c)$, the map $X_1(\lambda_0)$ on Λ , which is the displacement in x along the basic arc with $\lambda(0) = \lambda_0$ in the basic region, is a strictly increasing function and its range is $(0, \infty)$.

E. Duration of Basic Arcs and Idling Arcs

We compute the duration of basic arcs. Given an extremal $(x(t), y(t), \lambda(t))$ with $(y(0), \lambda(0)) = (0, \lambda_0) \in 0 \times \Lambda$, let T_1 be the smallest $t > 0$ such that $(y(t), \lambda(t)) = (\lambda_0, 0)$ where Λ is defined in (13) and (14). We can regard T_1 as a function of $\lambda_0 \in \Lambda$. Recall the three logic sentences, P_1 , P_2 and Q defined in Section II. We will compute T_1 separately for the following two cases:

$$[P_1 \vee Q] \quad \text{and} \quad P_2.$$

We begin with the case of $[P_1 \vee Q]$. The (y, λ) -projection of a basic arc is like the arc $ABCD$ in Fig. 7(a) and (b). It consists of the three arcs, AB , BC , and CD where $u = (y + \lambda)/(-2\alpha)$ on AB and CD and $u = +1$ on BC . Let us compute the flight time Δt_{AB} from A to B . By the S_1 symmetry, the flight time Δt_{CD} is the same as Δt_{AB} . The dynamics are given

in (48)–(50) with the initial condition $(y(0), \lambda(0)) = (0, \lambda_0)$. The solution is given by

$$\begin{cases} y(t) = \frac{\lambda_0}{-2\alpha\omega} \sin(\omega t) \\ \lambda(t) = \lambda_0 \cos(\omega t) + \frac{\lambda_0(1+2\alpha c)}{2\alpha\omega} \sin(\omega t). \end{cases} \quad (69)$$

where

$$\omega = \sqrt{\det A} = \sqrt{c(1+\alpha c)/(-\alpha)} > 0. \quad (70)$$

Thus

$$\Delta t_{AB}(\lambda_0) = \Delta t_{CD}(\lambda_0) = \frac{1}{\omega} \sin^{-1} \left(\frac{-2\alpha\omega y_B(\lambda_0)}{\lambda_0} \right) \quad (71)$$

where $y_B(\lambda_0)$ is given in (60). We now compute the flight time Δt_{BC} for BC . For this purpose, the corresponding y -dynamics can be written as

$$\dot{y} = -cy + 1 \quad y(0) = y_B \quad y(\Delta t_{BC}) = y_C$$

where y_B and y_C are given in (60) and (61). Direct integration yields

$$\Delta t_{BC}(\lambda_0) = \frac{1}{c} \ln \left(\frac{1 - cy_B(\lambda_0)}{1 - cy_C(\lambda_0)} \right). \quad (72)$$

By (71) and (72), the total flight time T_1 of $ABCD$ is given by

$$T_1(\lambda_0) = \frac{2}{\omega} \sin^{-1} \left(\frac{-2\alpha\omega y_B(\lambda_0)}{\lambda_0} \right) + \frac{1}{c} \ln \left(\frac{1 - cy_B(\lambda_0)}{1 - cy_C(\lambda_0)} \right). \quad (73)$$

We now consider the case of P2. The y -dynamics is

$$\dot{y} = -cy + 1 \quad y(0) = 0 \quad y(T_1) = \lambda_0.$$

Direct integration yields

$$T_1(\lambda_0) = \frac{1}{c} \ln \left(\frac{1}{1 - c\lambda_0} \right). \quad (74)$$

We have so far verified the formula of T_1 in (16).

We make a remark on the relation between X_1 and T_1 . Recall that $X_1(\lambda_0)$ is a bijection from Λ to $(0, \infty)$. The corresponding flight time $T_1(\lambda_0)$ can be regarded as function of $X_1 \in (0, \infty)$ as follows:

$$T_1(X_1) = T_1 \circ \lambda_0(X_1). \quad (75)$$

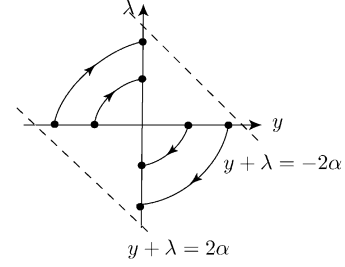


Fig. 8. Idling arcs do not contribute to the net displacement of a particle.

Recall the definition of *idling arcs* in Section II. An idling arc is an arc whose (y, λ) -projection starts from the positive (respectively, negative) y -axis, goes through the fourth quadrant (respectively, the second quadrant) in the region of R_I , and ends at the negative (respectively, positive) λ -axis; see Fig. 8. By the S_2 symmetry, they do not contribute to the displacement of x . Idling arcs occur in the case of $[P1 \vee Q]$. The flight time, T_{idling} , which we call *idling time*, of idling arcs is given by

$$T_{\text{idling}} = \frac{1}{\omega} \sin^{-1}(-2\alpha\omega) \quad (76)$$

with ω in (70). Notice that the idling time is independent of the coordinates of the initial point on the y -axis.

F. Construction and Uniqueness of Optimal Trajectories

We investigate the issue of the construction and uniqueness of the time-optimal trajectory. Recall that the (y, λ) projection of optimal trajectories must start from the λ -axis and end at the y -axis as $y(0) = 0$ and $\lambda(t_f) = 0$. Because of the S_3 symmetry, if a trajectory with $\lambda(0) \in \Lambda$ is optimal, then its S_3 image is also optimal. Hence, without loss of generality, we will always give proofs only for optimal trajectories starting with $\lambda(0) \in \Lambda$.

We call an extremal an n -shot extremal where $n \in \mathbb{N}$ if the maximum number of basic arcs in the extremal is n . Equivalently, we say that an extremal is an n -shot extremal if its (y, λ) projection meets with the y -axis n times. For example, the extremal whose (y, λ) projection is in Fig. 9(a), is a one-shot extremal and that in Fig. 9(b) is a two-shot extremal. The one in Fig. 9(c) corresponds to a multishot extremal. Notice that the (y, λ) projection of an n -shot extremal with $n \geq 3$ is a closed curve by symmetry. By the discussion in Sections VI-C and VI-D, we know that for a given $x_f > x_0$ there always exists a unique n -shot extremal for each $n \in \mathbb{N}$ with $\lambda(0) = \lambda_0 \in \Lambda$ reaching x_f . Hence, we need to know how to find time-optimal ones among them. We divide our discussion into the following two cases.

$$(3 + 4\alpha c) \geq 0 \quad (3 + 4\alpha c) < 0.$$

Case $(3 + 4\alpha c) \geq 0$: We will show that the two one-shot extremals are the time-optimal ones for a given $x_f > x_0$. Recall

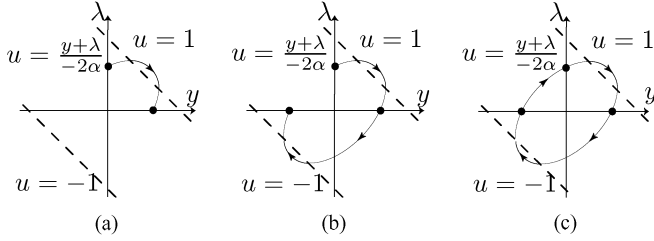


Fig. 9. There are several types of extremals. Their (y, λ) -projection must start from the λ -axis and end on the y -axis because of the boundary condition. (a) One-shot extremal, (b) two-shot extremal, and (c) multishot extremal.

the three logic sentences, P1, P2, and R defined in Section II. It is easy to verify

$$P1 \vee P2 \Rightarrow R.$$

We will divide our discussion into the following three cases:

$$P2, P1, [R \wedge \neg(P1 \vee P2)].$$

We first consider the case of P2. Suppose that the two-shot extremal, $FGHIJK$, in Fig. 10(a) is time-optimal with a final time $t_f > 0$ and a control u for $x_f > x_0$. Let us consider another control $v(t) = 1$ for $t \in [0, t_f]$. Let $FGHJ'K'$ be the trajectory due to v . By the S_2 symmetry, $\Delta x_{HI} = 0$. Hence, the increment in x due to the control u is given by

$$\Delta x_{FGHIJK}(u) = \Delta x_{FGH}(u=1) + \Delta x_{S_3(I)FG}(u=1)$$

by the S_3 symmetry where $S_3(IJK) = S_3(I)FG$. The increment in x due to the control v is given by

$$\Delta x_{FGHJ'K'}(v) = \Delta x_{FGH}(v=1) + \Delta x_{HJ'K'}(v=1).$$

Along $S_3(I)FGHJ'K'$, $\dot{y} = -cy + 1 > 0$. Let J' be the point such that the flight time on $S_3(I)FG$ is the same as that on HJ' . Namely, $\Delta t_{S_3(I)FG} = \Delta t_{HJ'}$. Then the flight time on HI is the same as that on $J'K'$. Notice that

$$\begin{aligned} \Delta x_{S_3(I)FG}(u=1) &= \int_0^{\Delta t_{S_3(I)FG}} (y_1(t) + \alpha) dt \\ &< \int_0^{\Delta t_{HJ'}} (y_2(t) + \alpha) dt \\ &= \Delta x_{HJ'}(v=1) \end{aligned}$$

where $y_1(t) = y_{S_3(I)}e^{-ct} + (1)/(c)(1 - e^{-ct})$, $y_2(t) = y_H e^{-ct} + (1)/(c)(1 - e^{-ct})$, and $y_{S_3(I)} < y_H$. On $J'K'$, $\dot{x} = y + \alpha > 0$ because $y_H \geq y_G = \lambda_0 \geq -2\alpha$ and $\dot{y} > 0$ on $J'K'$. Thus, $\Delta x_{J'K'}(v=1) > 0$. Therefore, $\Delta x_{FGHJ'K'}(v) > \Delta x_{FGHIJK}(u)$. This implies, by the continuity of $t \mapsto x(t) \in \mathbb{R}$, that the trajectory due to the control

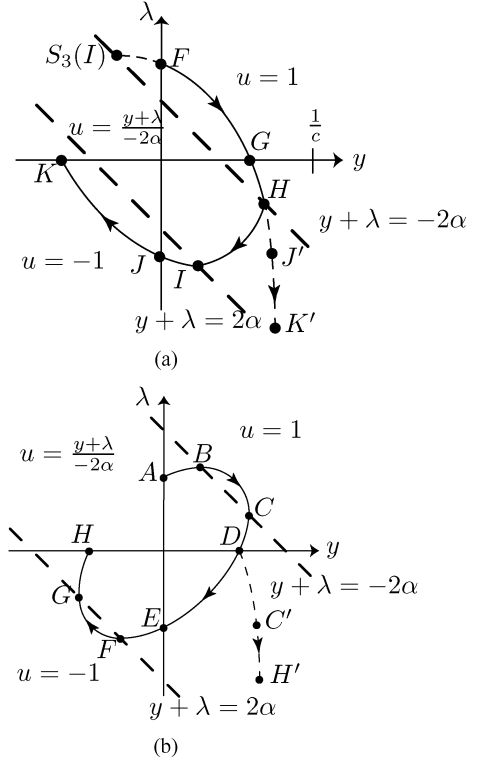


Fig. 10. One-shot extremal is the optimal one among extremals for a given $x_f > x_0$ when $(1 + 2\alpha c) > 0$. (a) $\lambda_0 \in \Lambda_2 = [-2\alpha, 1/c]$. (b) $\lambda_0 \in \Lambda_1 = (-2\alpha\sqrt{1+\alpha c}, -2\alpha)$.

v reaches x_f before $t = t_f$, which contradicts the time-optimality of the control u . Hence, no two-shot extremals can be time-optimal in the case of P2. Similarly, one can show that no multi-shot extremals can be time-optimal in the case of P2.

We now consider the case of P1. Suppose that a two-shot extremal with $\lambda_0 \in \Lambda_1$ is time-optimal with $t_f > 0$ and control u for $x_f > x_0$. Let $ABCDEFGH$ be the trajectory associated with u in Fig. 10(b). Let t_1 be the time when the trajectory reaches D . Let us construct another control v as follows:

$$v(t) = \begin{cases} u(t), & \text{if } t \in [0, t_1] \\ 1, & \text{if } t \in [t_1, t_f]. \end{cases}$$

Let $ABCD C'H'$ be the trajectory associated with the control v where $DC'H'$ corresponds to the time interval $[t_1, t_f]$, and we choose C' such that $\Delta t_{BC} = \Delta t_{DC'}$. By the S_1 and S_2 symmetry

$$\Delta x_{AB} + \Delta x_{CD} = \Delta x_{EF} + \Delta x_{GH} = 0.$$

The total increment in x due to the control u is given by

$$\Delta x(u) = 2\Delta x_{BC}. \quad (77)$$

The total increment in x due to the control v is

$$\Delta x(v) = \Delta x_{BC} + \Delta x_{DC'} + \Delta x_{C'H'}. \quad (78)$$

By the S_1 symmetry, $y_D = \lambda_0$, so

$$y_D = \lambda_0 > -2\alpha\sqrt{1+\alpha c} \geq -\alpha \geq y_B \quad (79)$$

where $\lambda_0 \in \Lambda$, $(3 + 4\alpha c) \geq 0$ and (60) were used for the three inequalities. As $\dot{y} = -cy + 1 > 0$ on BC and DC' , we have

$$\begin{aligned} \Delta x_{BC} &= \int_0^{\Delta t_{BC}} (y_1(t) + \alpha) dt \\ &< \int_0^{\Delta t_{DC'}} (y_2(t) + \alpha) dt = \Delta x_{DC'} \end{aligned} \quad (80)$$

where $y_1(t) = y_B e^{-ct} + (1)/(c)(1 - e^{-ct})$ and

$$y_2(t) = y_D e^{-ct} + \frac{1}{c}(1 - e^{-ct}).$$

It is straightforward to see

$$\Delta x_{C'H'} > 0. \quad (81)$$

By (77), (78), (80), and (81),

$$\Delta x(v) > \Delta x(u).$$

This implies that the trajectory due to the control v reaches x_f before $t = t_f$, which contradicts the time-optimality of u . Therefore, there are no two-shot time-optimal trajectories in the case of P1. In the similar manner, one can show that no multishot extremals can be time-optimal in the case of P1.

We now consider the case of $[R \wedge \neg(P1 \vee P2)]$. In this case, all the basic arcs are of form $\Gamma_{\pm}^{\lambda_0}$ in (17). Notice that for the case of P1 we only used the fact that the basic arcs are of form $\Gamma_{\pm}^{\lambda_0}$ and $(3 + 4\alpha c) \geq 0$ where the latter was used in (79). Hence, the case of $[R \wedge \neg(P1 \vee P2)]$ can be handled in the same way as the case of P1, to conclude that no multishot extremals are time-optimal.

So far, we have showed that multishot extremals cannot be time-optimal when $(3 + 4\alpha c) \geq 0$. By the discussion in Section VI-D, we know that there exists a unique one-shot extremal with $\lambda_0 \in \Lambda$, or a basic arc, for a given $x_f > x_0$ such that $x(t_f) = x_f$. Hence, this basic arc and its image under the reflection S_3 are the only optimal trajectories.

Claim 6.4: If $(3 + 4\alpha c) \geq 0$, then there are exactly two time optimal trajectories for $x_f > x_0$. One is the one-shot extremal $\Gamma_{+}^{\lambda_0}$ (or, the basic arc) with $\lambda(0) = \lambda_0 = X_1^{-1}(x_f - x_0)$ and the other is its image by the reflection S_3 (see also Section II). The corresponding minimum time is $T_1(\lambda_0)$. The maps X_1 and T_1 are defined in (15) and (16).

Case $(3 + 4\alpha c) < 0$: In this case, unlike the case of $(3 + 4\alpha c) \geq 0$, we have no general proof that only one-shot extremals are time-optimal. Instead we provide a finite algorithm of finding all time-optimal trajectories for $x_f > x_0$.

Take $x_f > x_0$. For each $n \in \mathbb{N}$, there exists a unique n -shot extremal with $\lambda(0) = \lambda_{0,n} := X_1^{-1}((x_f - x_0)/n)$ reaching x_f . Since $(3 + 4\alpha c) < 0$ implies $(1 + 2\alpha c) \leq 0$, this n -shot extremal consists of n one-shot extremals and $(n-1)$ idling arcs between the n one-shot extremals; see (19), Figs. 7(a), and 9(c).

Each of the n one-shot extremals contributes $(x_f - x_0)/n$ to the increment in x , and each idling arc makes zero contribution. By this decomposition, the total time cost T_n for this n -shot extremal is

$$T_n = n \cdot T_1 \circ X_1^{-1}((x_f - x_0)/n) + (n-1) \cdot T_{\text{idling}} \quad (82)$$

where $T_1 \circ X_1^{-1}((x_f - x_0)/n)$ is the flight time corresponding to the increment $(x_f - x_0)/n$ in x where X_1, T_1 , and T_{idling} are given in (63), (73), and (76). There exists $N \in \mathbb{N}$ such that

$$T_1 = T_1 \circ X_1^{-1}(x_f - x_0) < N \cdot T_{\text{idling}}$$

which implies

$$T_1 < T_n \quad \forall n \geq (N+1).$$

Thus, all possible time-optimal trajectories are among the first N extremals. Choose $k \in \{1, \dots, N\}$ such that $T_k = \min\{T_i : 1 \leq i \leq N\}$. Then, the k -shot extremal is a time-optimal trajectory. Such k 's give all the time-optimal trajectories corresponding to $x_f > x_0$. This proves the procedure [B.1]–[B.4] in Section II, which can be summarized as follows.

Claim 6.5: If $(3 + 4\alpha c) < 0$, then there is a finite and explicit procedure ([B.1]–[B.4] in Section II) of finding all time-optimal trajectories and the corresponding minimum time for $x_f > x_0$.

As a remark, we give a practical way of showing that the one-shot extremals are the unique time-optimal trajectories. First, with (63) and (73), one draws the graph $(X_1, T_1(X_1))$ in (75). Suppose that it is strictly concave. Then

$$T_1 = T_1(X_1 = \Delta x) < n \cdot T_1(X_1 = \Delta x/n) \quad \forall n = 2, 3, \dots$$

which, with (82), implies that $T_1 < T_n$ for all $n \geq 2$. Hence, the two one-shot extremals are the only time-optimal trajectories if $T_1(X_1)$ is a strictly concave function of X_1 .

G. Initial Undershoots

One can check that for $\lambda(0) \in \Lambda$, every extremal has $\dot{x}(t) \leq 0$ from $t = 0$ until $y(t)$, which is initially zero, becomes $-\alpha > 0$. In other words, $x(t)$ goes through an initial undershoot. We will compute this undershoot and its duration. Here, we do not give detailed computation because the methodology is very similar to those in Sections VI-D and VI-E. Recall the three logic sentences, P1, P2, and Q defined in Section II.

Let us first consider the case of $[P1 \vee Q]$. The undershoot consists of two parts in Fig. 7(a) and (b): AB and $B \rightarrow (-\alpha, \bar{\lambda})$ for some $\bar{\lambda}$ where $(-\alpha, \bar{\lambda})$ lies between B and C on the trajectory. One can compute the duration $T_{\text{under}}(\lambda_0)$ of the undershoot as

$$T_{\text{under}}(\lambda_0) = \Delta t_{AB}(\lambda_0) + \frac{1}{c} \ln \left(\frac{1 - cy_B(\lambda_0)}{1 + c\alpha} \right)$$

and the (negative-valued) amount of the undershoot

$$X_{\text{under}}(\lambda_0) = x(T_{\text{under}}(\lambda_0)) - x(0)$$

as

$$\begin{aligned}
& X_{\text{under}}(\lambda_0) \\
&= \frac{\lambda_0^2}{8\omega^3\alpha^2} \left\{ (\omega^3\alpha + \omega c + \omega\alpha c^2)t \right. \\
&\quad + (1 + 2\alpha c)\omega \sin^2(\omega t) \\
&\quad + (\omega^2\alpha - c - \alpha c^2) \cos(\omega t) \sin(\omega t) \left. \right\} |_{t=\Delta t_{AB}(\lambda_0)} \\
&\quad + \frac{1}{c} \left\{ y_B(\lambda_0) + \alpha + \left(\alpha + \frac{1}{c} \right) \right. \\
&\quad \times \ln \left(\frac{1 - cy_B(\lambda_0)}{1 + c\alpha} \right) \left. \right\}
\end{aligned}$$

where ω , Δt_{AB} , and $y_B(\lambda_0)$ are defined in (70), (71), and (60).

In the case of P2, the undershoot corresponds to the arc, $F \rightarrow (-\alpha, \bar{\lambda})$ for some $\bar{\lambda}$ in Fig. 7(b) where $(-\alpha, \bar{\lambda})$ lies between F and G . One can compute the duration $T_{\text{under}}(\lambda_0)$ of the undershoot as

$$T = \frac{1}{c} \ln \left(\frac{1}{1 + \alpha c} \right)$$

and the amount $X_{\text{under}}(\lambda_0)$ of the undershoot as

$$X_{\text{under}}(\lambda_0) = \frac{1}{c} \left(\alpha + \left(\alpha + \frac{1}{c} \right) \ln \left(\frac{1}{1 + \alpha c} \right) \right).$$

Notice that this undershoot is independent of λ_0 .

H. Discussion on State Constraints on x

We have seen in Section III that physics imposes a constraint on x as either

$$x \geq 0 \quad \text{or} \quad x \leq 0. \quad (83)$$

We have not considered these constraints in computing optimal trajectories in this paper. We now make a remark that in some cases optimal trajectories derived without the state constraints are also optimal with the state constraints.

First, let us consider the result in Section IV again—this time with the state constraints in (83). In either case of (83), it is not hard to verify that the result in Section IV still holds without violating the constraint. Namely, there are no time optimal trajectories for $x_f < x_0$ even in the existence of either state constraint in (83).

Second, we consider the constraint $x \leq 0$ with the initial and final condition satisfying $x_f \geq x_0$. It is trivial to see that an optimal trajectory derived without the state constraint reaches x_f without violating the constraint, so it is still optimal in the existence of the state constraint.

Lastly, we consider the constraint $x \geq 0$ with the initial and final condition satisfying $x_f \geq x_0$. We divide the discussion into the two cases; $(3 + 4\alpha c) \geq 0$ and $(3 + 4\alpha c) < 0$. When $(3 + 4\alpha c) \geq 0$, we know from Section VI-F that there exists a unique $\lambda_0 \in \Lambda$ such that the two basic arcs $\Gamma_{\pm}^{\lambda_0}$ are the only optimal trajectories. Recall the discussion on initial undershoots in Section VI.G. Also, recall from Section VI-D that X_1 is a positive function. We have $\dot{x}(t) > 0$ after $x(t)$

passes through $x_0 + X_{\text{under}}(\lambda_0)$ along the basic arc. Hence, we have $\min_{t \in [0, T_1(\lambda_0)]} x(t) = x_0 + X_{\text{under}}(\lambda_0)$ where $T_1(\lambda_0)$ is the duration of the basic arc. It follows that in the case of $(3 + 4\alpha c) \geq 0$ if $x_0 \geq |X_{\text{under}}(\lambda_0)|$ then the optimal trajectories derived without the state constraint $x \geq 0$ are optimal with the state constraint. On the other hand, as seen in Section VI-F, when $(3 + 4\alpha c) < 0$ there can be more than two optimal trajectories for a given $x_f \geq x_0$. Some of the optimal trajectories can be multishot extremals. On an idling arc, one has $\dot{x} = (\lambda^2 - y^2)/(-4\alpha)$. Hence, during the first half of an idling arc $\dot{x}(t) \geq 0$ and during the second half $\dot{x}(t) \leq 0$. Recall that an idling arc does not contribute any net displacement in x . Let t_1 and $t_2 = t_1 + T_{\text{idling}}$ be the initial and final time of an idling arc of an optimal trajectory with $\lambda_0 \in \Lambda$ derived without the state constraint $x \geq 0$. Then, it follows that $\min_{t \in [t_1, t_2]} x(t) = x(t_1) = x(t_2) \geq x_0$. Considering the decomposition of multishot extremals in (19), one can see that in the case of $(3 + 4\alpha c) < 0$ if $x_0 \geq |X_{\text{under}}(\lambda_0)|$ then the optimal trajectories derived without the state constraint $x \geq 0$ remain optimal with the state constraint. Therefore, irrespective of the sign of $(3 + 4\alpha c)$, an optimal trajectory with $\lambda_0 \in \Lambda$ derived without the state constraint $x \geq 0$ is again optimal with the state constraint if $x_f > x_0 \geq |X_{\text{under}}(\lambda_0)|$.

VII. SYMPLECTIC REDUCTION PICTURE

We cast the analysis used in Sections V and VI in the framework of symplectic reduction theory in order to provide geometric insight into the problem. We refer readers to [1] and [9] for symplectic reduction theory.

Let $G = \mathbb{R}$ be the Lie group acting on $M = \mathbb{R}^2 = \{(x, y)\}$ by translation in x and on the cotangent bundle $T^*M = \mathbb{R}^4 = \{(x, y, \lambda_x, \lambda_y)\}$ by cotangent lift where T^*M is equipped with the canonical symplectic form, $\Omega = dx \wedge d\lambda_x + dy \wedge d\lambda_y$. The momentum map $J : T^*M \rightarrow \mathbb{R}$ corresponding to the G -action is $J(x, y, \lambda_x, \lambda_y) = \lambda_x$. We have $J^{-1}(\nu) \simeq \mathbb{R}^3$ for any $\nu \in \mathbb{R}$, and $J^{-1}(\nu)/G \simeq \mathbb{R}^2$. One can construct the symplectic projection $\pi : J^{-1}(\nu) \subset T^*M \rightarrow J^{-1}(\nu)/G \simeq \mathbb{R}^2$ by $\pi(x, y, \nu, \lambda_y) = (y, \lambda_y)$ where $J^{-1}(\nu)/G$ is given the canonical symplectic form $\omega = dy \wedge d\lambda_y$. Notice that the Hamiltonian H in (32) is G -invariant, that the control u in (38) is G -invariant, and that the Hamiltonian M° in (33) is also G -invariant. Hence, H and M° induce reduced Hamiltonians h and m° on $J^{-1}(\nu)/G$ as follows:

$$\begin{aligned}
h(y, \lambda_y, u) &= H(x, y, \nu, \lambda_y, u) \\
&= \lambda_x \alpha u^2 + (\lambda_x y + \lambda_y)u - cy\lambda_y \quad (84)
\end{aligned}$$

$$m^\circ(y, \lambda_y) = M^\circ(x, y, \nu, \lambda_y) = \max_{|u| \leq 1} h(y, \lambda_y, u). \quad (85)$$

Notice that the control u maximizing $H(x, y, \nu, \lambda_y, u)$ is the same as the control maximizing $h(y, \lambda_y, u)$ because the control u maximizing $H(x, y, \nu, \lambda_y, u)$ is G -invariant and the group G is abelian. Equations (35) and (37) are the (reduced) dynamics of the reduced Hamiltonian m° (or h) on the based space $(J^{-1}(\nu)/G, \omega)$, and (34) can be regarded as a reconstruction equation to compute the displacement along the fiber $G = \mathbb{R}$. The map X_1 constructed in Section VI.D measures this displacement while $(y(t), \lambda(t))$ with $(y(0), \lambda(0)) \in 0 \times \Lambda$ stays in

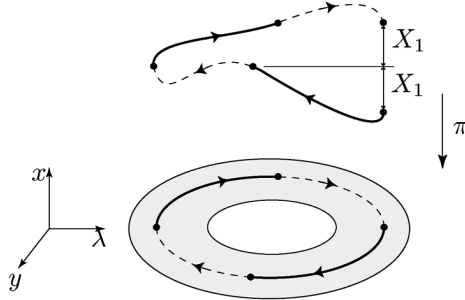


Fig. 11. Level set $J^{-1}(\nu)$ is diffeomorphic to $\mathbb{R}^3 = \{(x, y, \lambda)\}$ where the x -axis is the fiber, and $\pi : J^{-1}(\nu) \rightarrow \mathbb{R}^2$ is the projection to the y - λ plane. The map X_1 measures the displacement along the fiber. The dotted arcs are the idling arcs. The shaded region here corresponds to those in Fig. 4.

the first quadrant in the y - λ plane. In particular, it is interesting to notice that if $[(1 + 2\alpha c) \leq 0]$ or $[(1 + 2\alpha c) > 0 \text{ and } \lambda(0) \in \Lambda_1]$, then X_1 measures the half of the total phase (or, holonomy with respect to the trivial connection on the principal bundle $\pi : J^{-1}(\nu) \rightarrow J^{-1}(\nu)/G$) corresponding to a closed trajectory $(y(t), \lambda(t))$ in the base space; see Fig. 11.

In Sections VI.A and VI.B, we did not use the reduced space (\mathbb{R}^2, ω) and the projection π constructed above when $J = \lambda_x = \nu \neq 0$, but we used the scaled variable $\lambda = \lambda_y/\lambda_x$ in (41) and (43) so as to show the discrete symmetry as reflection maps and to deal with all cases of $\lambda_x > 0$ simultaneously. To put this scaling in the reduction process, one can use $\nu dy \wedge d\lambda$ as a symplectic form on the reduced space $J^{-1}(\nu)/G = \{(y, \lambda)\}$, and define the projection $\pi_\nu : J^{-1}(\nu) \rightarrow J^{-1}(\nu)/G$ by

$$\pi_\nu : (x, y, \nu, \lambda_y) \mapsto (y, \lambda) = (y, \lambda_y/\nu).$$

In this case, the reduced Hamiltonians h and m° induced by H and M° are those in (84) and (85) with λ_y replaced by $\nu\lambda$. The vector field $(\dot{y}, \dot{\lambda})$ in Section VI-B is the Hamiltonian vector field of the reduced Hamiltonian m° (or h) on the reduced space $(\mathbb{R}^2, \nu dy \wedge d\lambda)$.

Remark 7.1: The $G = \mathbb{R}$ symmetry in the Hamiltonian dynamics was created by the nonlinear coordinate change in (23). This reduction process extends to the case where we consider many particles on the one-dimensional invariant line.

VIII. SIMULATIONS

We demonstrate the theoretical result by simulation. Suppose that the particle and the medium are given such that

$$\alpha = -3/4 \quad c = 1$$

and that the initial and final position of the particle is given by

$$x_0 = 1 \quad x_f = 2.$$

These numbers are chosen arbitrarily, but one can also do the same simulation with real data once they are given. Let $\Delta x = x_f - x_0$. One can check that the condition $(1 + \alpha c) > 0$ for the

TABLE I
TIME COSTS OF THE FIRST FIVE EXTREMALS

n	$\lambda_{0,n}$	$T_1 \circ X_1^{-1}(\Delta x/n)$	Time cost T_n
1	0.8649	7.8117	7.8117
2	0.8577	5.7950	13.4038
3	0.8491	5.1093	18.9556
4	0.8415	4.7597	24.4802
5	0.8350	4.5463	29.9867

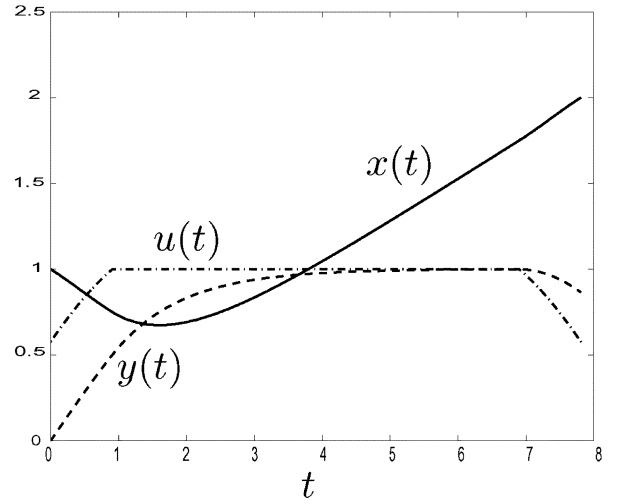


Fig. 12. State $(x(t), y(t))$ of the optimal trajectory and the corresponding optimal control $u(t)$ for $x_0 = 1$ and $x_f = 2$ with minimum time $t_f = 7.8117$.

existence of optimal trajectories holds. The idling time in (76) is given by $T_{\text{idling}} = 1.8138$. By (63) and (71)–(73), the flight time of the one-shot extremal T_1 is given by $T_1 = 7.8117$. By the algorithm in Claim 6.5, one learns that the minimum time is in $\{T_1, \dots, T_5\}$. By (82) one computes these flight times in Table I.

Hence, the one-shot extremal is the optimal trajectory and the minimum time is $t_f = 7.8117$. This result agrees with the uniqueness result because the given α and c satisfy $(3 + 4\alpha c) \geq 0$. The optimal control $u(t)$ and the optimal trajectory $(x(t), y(t))$ are given in Fig. 12. Notice the initial undershoot predicted in Section VI-G.

IX. CONCLUSION AND FUTURE WORK

The time-optimal control problem for the dielectrophoretic system studied in this paper has several interesting features. The existence of a term quadratic in control creates the non-existence of optimal trajectories when the final position of the particle is below the initial position. In contrast, when the final position is above the initial position we can show the existence and uniqueness of optimal trajectories in a range of the parameters α and c . In the other range of α and c we give a finite algorithm of finding all optimal trajectories instead. Both continuous and discrete symmetry in the problem simplifies the analysis.

In summary, the optimal trajectories are described as follows. There are three different types of optimal trajectories depending on the values of α and c , and the displacement of the particle. If $(1 + 2\alpha c) > 0$, basic arcs, or one-shot extremals, are the

optimal trajectories. To move a particle a long distance, the optimal control is the saturated controls, i.e., $u = \pm 1$, but for a small displacement of the particle the optimal control consists of three parts; a linear control, a saturated control and a linear control again. If $(1 + 2\alpha c) \leq 0$ and $(3 + 4\alpha c) \geq 0$, then for any displacement of the particle the optimal control consists of three parts; a linear control, a saturated control and a linear control again. If $(3 + 4\alpha c) < 0$, optimal trajectories may have the structure of multi-shot extremals in (18) and (19).

As for future work, we will take into account a state constraint, and/or consider the time-optimal control of two different particles for the purpose of separating them which are initially close to each other. They have important applications in nano/bio-technology [5], [7], [8]. We believe that the use of control systems theory will refine and improve the manipulation of particles in applications and that our work in this paper makes a first forward step in this direction.

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